



PITP, UBC, Vancouver 2019

**Quantum Controlling
Levitated Solids:
*a novel probe for the
gravity-quantum interface***

Markus Aspelmeyer

Vienna Center for Quantum Science and Technology (VCQ)

Faculty of Physics, University of Vienna, Austria

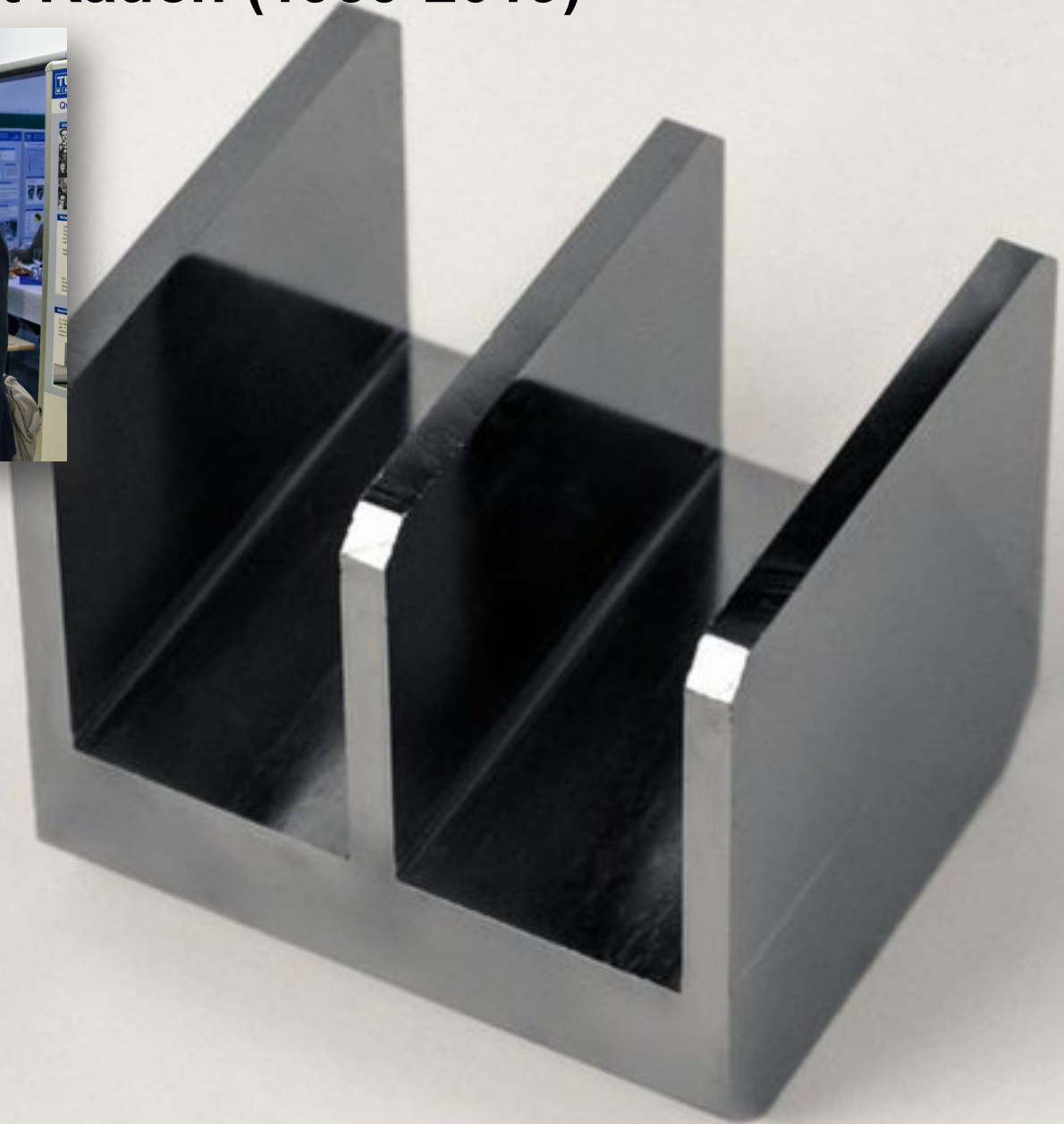
IQOQI, Austrian Academy of Sciences

In memory of Helmut Rauch (1939-2019)



The neutron interferometer

*Rauch, Treimer, Bonse,
Physics Letters A 47, 369 (1974)*



Observation of Gravitationally Induced Quantum Interference*

R. Colella and A. W. Overhauser

Department of Physics, Purdue University, West Lafayette, Indiana 47907

and

S. A. Werner

Scientific Research Staff, Ford Motor Company, Dearborn, Michigan 48121

(Received 14 April 1975)

We have used a neutron interferometer to observe the quantum-mechanical phase shift of neutrons caused by their interaction with Earth's gravitational field.

$$\Delta\gamma = \frac{1}{\hbar} \int m \underbrace{\Delta\phi}_{\substack{\downarrow \\ \text{gravitational potential} \\ (\text{on Earth: } \phi = g \cdot h)} dt$$

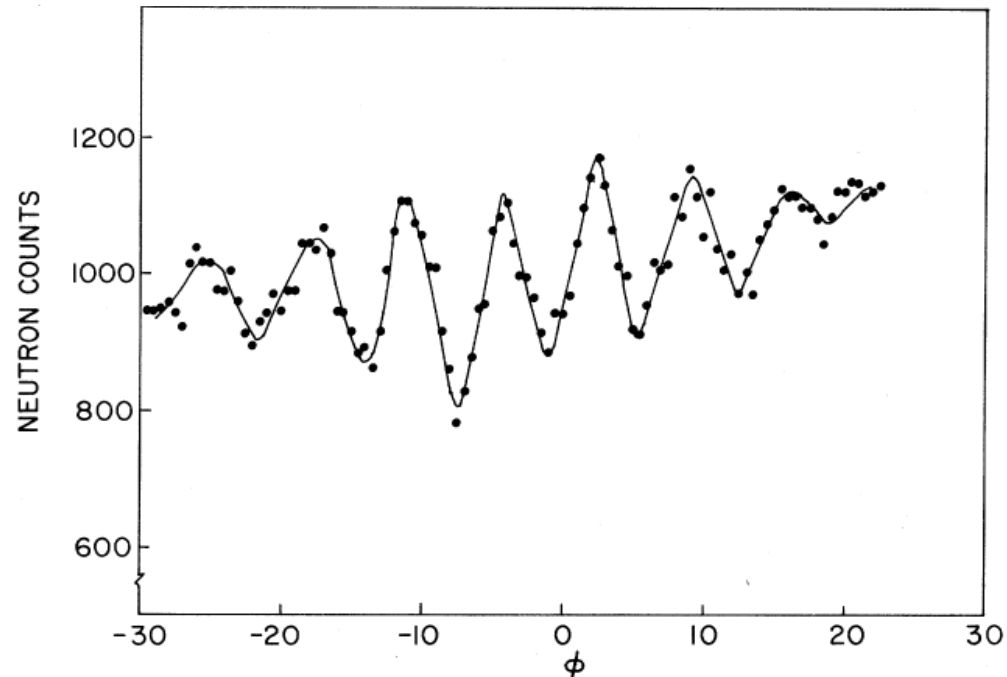
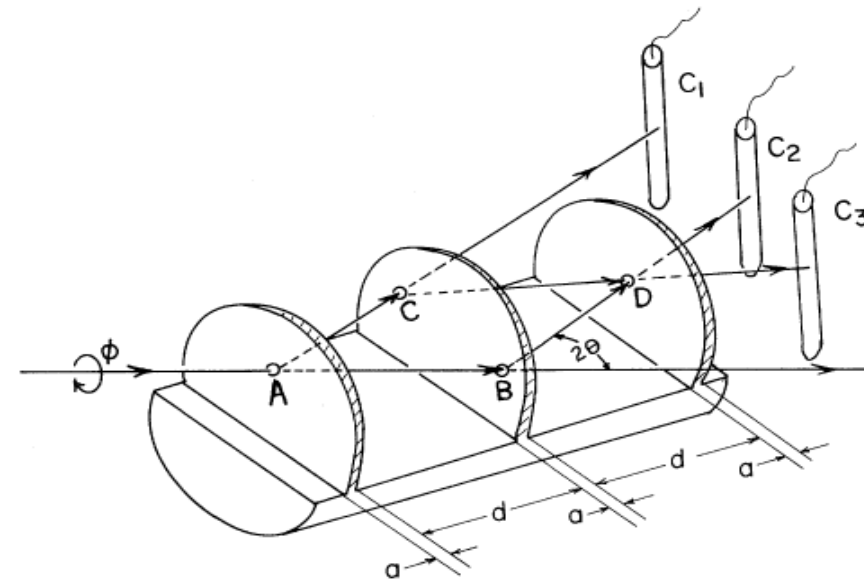


FIG. 1. Schematic diagram of the neutron interferometer and ^3He detectors used in this experiment.

Observation of Gravitationally Induced Quantum Interference*

R. Colella and A. W. Overhauser

Department of Physics, Purdue University, West Lafayette, Indiana 47907

and

S. A. Werner

Scientific Research Staff, Ford Motor Company, Dearborn, Michigan 48121

(Received 14 April 1975)

We have used a neutron interferometer to observe the quantum-mechanical phase shift of neutrons caused by their interaction with Earth's gravitational field.

$$\Delta\gamma = \frac{1}{\hbar} \int m \underbrace{\Delta\phi}_{\substack{\downarrow \\ \text{gravitational potential} \\ (\text{on Earth: } \phi = g \cdot h)} dt$$

Nature 1999

Measurement of gravitational acceleration by dropping atoms

Achim Peters, Keng Yeow Chung & Steven Chu

Physics Department, Stanford University, Stanford, California 94305-4060, USA

Laser-cooling of atoms and atom-trapping are finding increasing application in many areas of science¹. One important use of laser-cooled atoms is in atom interferometers². In these devices, an atom is placed into a superposition of two or more spatially separated atomic states; these states are each described by a quantum-mechanical phase term, which will interfere with one another if they are brought back together at a later time. Atom

Nature 2002

Quantum states of neutrons in the Earth's gravitational field

Valery V. Nesvizhevsky*, Hans G. Börner*, Alexander K. Petukhov*, Hartmut Abele†, Stefan Baeßler†, Frank J. Rueß†, Thilo Stöferle†, Alexander Westphal†, Alexei M. Gagarski‡, Guennady A. Petrov‡ & Alexander V. Strelkov§

* *Institute Laue-Langevin, 6 rue Jules Horowitz, Grenoble F-38042, France*

† *University of Heidelberg, 12 Philosophenweg, Heidelberg D-69120, Germany*

‡ *Petersburg Nuclear Physics Institute, Orlova Roscha, Gatchina, Leningrad reg. R-188350, Russia*

§ *Joint Institute for Nuclear Research, Dubna, Moscow reg. R-141980, Russia*

The discrete quantum properties of matter are manifest in a variety of phenomena. Any particle that is trapped in a sufficiently deep and wide potential well is settled in quantum bound states. For example, the existence of quantum states of electrons in an electromagnetic field is responsible for the structure of atoms¹⁶, and quantum states of nucleons in a strong nuclear field give rise to the structure of atomic nuclei¹⁷. In an analogous way, the gravitational field should lead to the formation of quantum states.

Observation of Gravitationally Induced Quantum Interference*

R. Colella and A. W. Overhauser

Purdue University, West Lafayette, Indiana 47907

(Kasevich group)

1991 $\Delta g/g = 1 \times 10^{-6}$

1998 $\Delta g/g = 3 \times 10^{-8}$

2014 $\Delta g/g = 5 \times 10^{-13}$

We h
of neut

Nature 1999

*Dearborn, Michigan 48121
1975)*

the quantum-mechanical phase shift
gravitational field.

Nature 2002

$$\Delta \gamma = \frac{1}{\hbar} \int m \Delta \phi dt$$

↓
gravitational potential
(on Earth: $\phi = g \cdot h$)

Measurement of gravitational acceleration by dropping atoms

Achim Peters, Keng Yeow Chung & Steven Chu

Physics Department, Stanford University, Stanford, California 94305-

Laser-cooling of atoms and atom-trapping are finding increasing application in many areas of science¹. One important use of laser-cooled atoms is in atom interferometers². In these devices, an atom is placed into a superposition of two or more spatially separated atomic states; these states are each described by a quantum-mechanical phase term, which will interfere with one another if they are brought back together at a later time. Atom

(Kasevich/Tino groups)

2007 $\Delta G/G = 3 \times 10^{-3}$

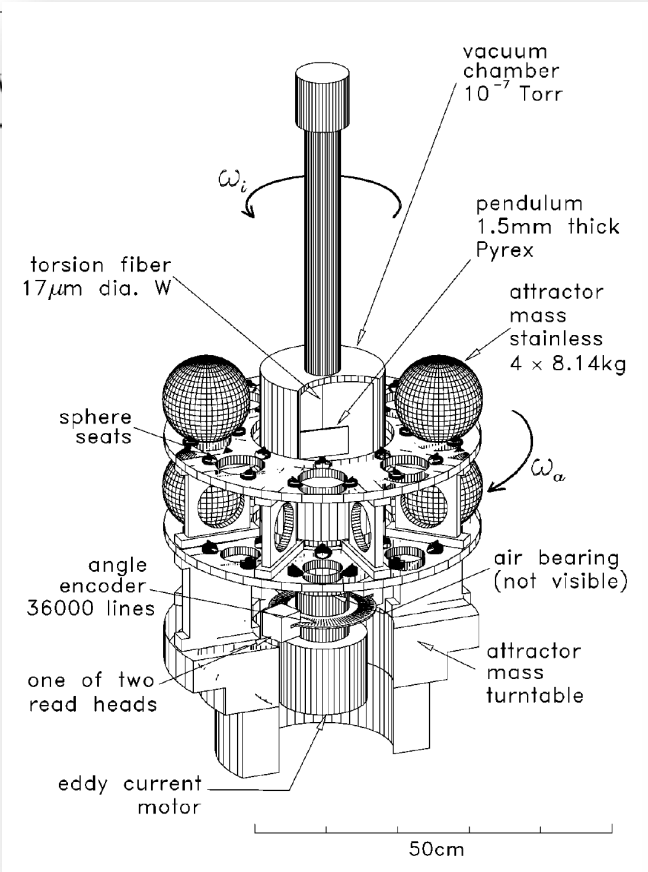
2014 $\Delta G/G = 1 \times 10^{-4}$

mainly limited by position of atoms

The discrete quantum properties of matter are manifest in a variety of phenomena. Any particle that is trapped in a sufficiently deep and wide potential well is settled in quantum bound states. For example, the existence of quantum states of electrons in an electromagnetic field is responsible for the structure of atoms¹⁶, and quantum states of nucleons in a strong nuclear field give rise to the structure of atomic nuclei¹⁷. In an analogous way, the gravitational field should lead to the formation of quantum states.

Quantum systems as sensitive probes of gravity

9 JUNE 1975



$$\Delta\gamma = \frac{1}{h} \int m \Delta\phi dt$$

↓
gravitational potential
(on Earth: $\phi = g \cdot h$)

Gundlach et al.,
PRL 2000
 $\Delta G/G = 15\text{ppm}$

$2014 \Delta G/G = 1 \times 10^{-4}$

mainly limited by position of atoms

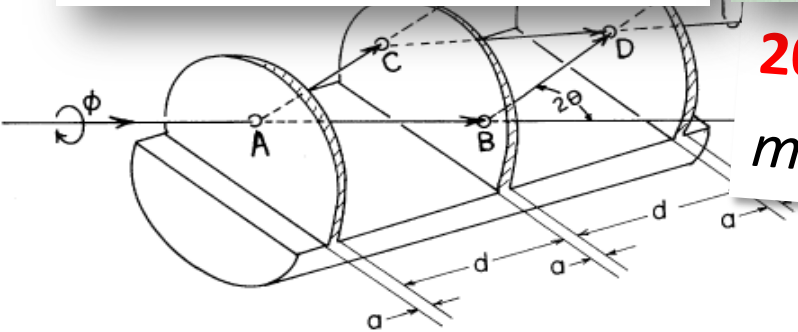
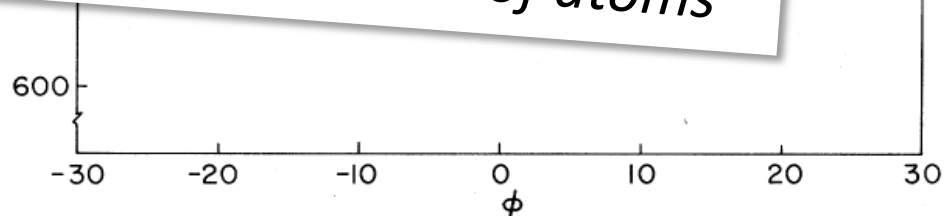
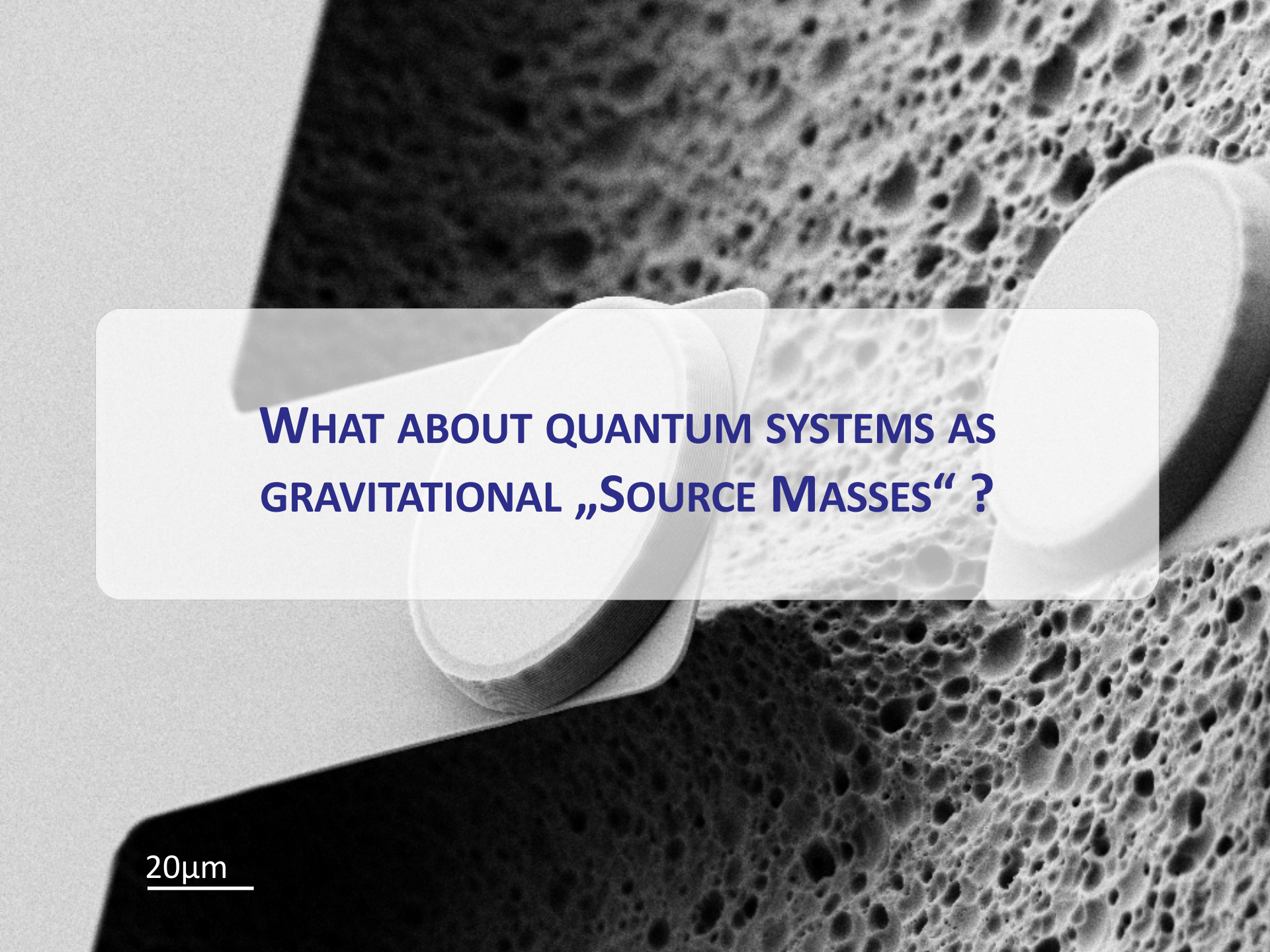


FIG. 1. Schematic diagram of the neutron interferometer and ^3He detectors used in this experiment.

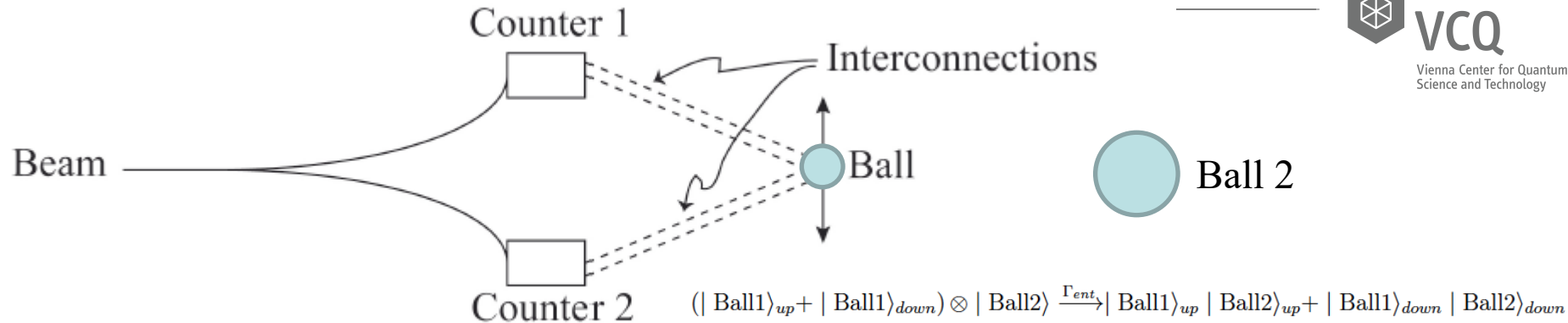


A scanning electron micrograph (SEM) showing a highly porous, sponge-like material with a complex, interconnected network of fibers. Two white, oval-shaped structures are visible, one in the foreground and one in the background, both appearing to be attached to or resting on the porous surface. The structures have a smooth, slightly curved top surface and a more textured, possibly layered or fibrous base. The overall image is in grayscale, highlighting the intricate texture of the porous material.

**WHAT ABOUT QUANTUM SYSTEMS AS
GRAVITATIONAL „SOURCE MASSES“ ?**

20 μ m

An ultimate experiment? Entanglement by gravity...

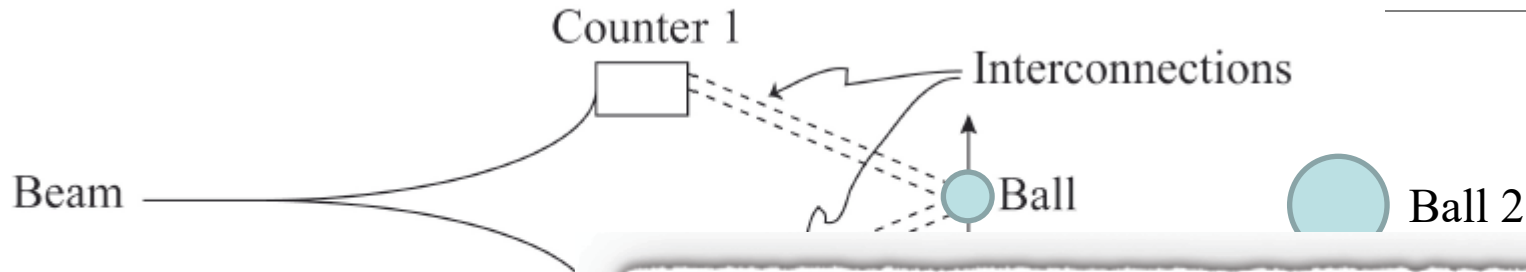


FEYNMAN: “Therefore, there must be an **amplitude for the gravitational field**, *provided* that the amplification necessary to reach a mass which can produce a gravitational field big enough to serve as a link in the chain does not destroy the possibility of keeping quantum mechanics all the way. There is a *bare* possibility (which I shouldn’t mention!) that quantum mechanics fails and becomes classical again when the amplification gets far enough, because of some minimum amplification which you can get across such a chain. But aside from that possibility, **if you believe in quantum mechanics up to any level then you have to believe in gravitational quantization in order to describe this experiment.**”

Chapel Hill Conference 1957 (29)

WITTEN: “What prevents this from becoming a practical experiment?”

An ultimate experiment? Entanglement by gravity...



Entangling via gravity

Entangling time:

$$t_{ent} = \frac{2\hbar}{GM^2} R^3 \left(\frac{1}{x_0}\right)^2$$

Pino et al., arXiv:1603.01553

$|12\rangle_{down}$

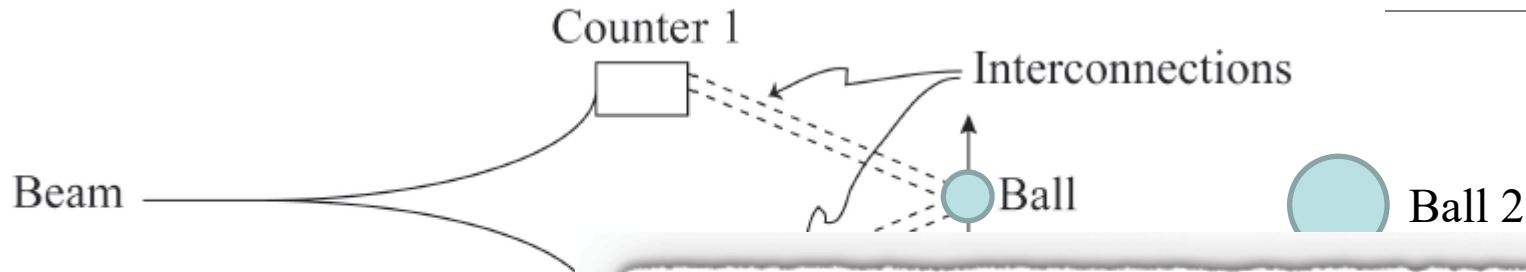
FEYNMAN: “Therefore, provided that the duce a gravitation destroy the possi is a bare possibilit fails and becomes because of some chain. But aside

ics up to any level then you have to believe in gravitational quantization in order to describe this experiment.”

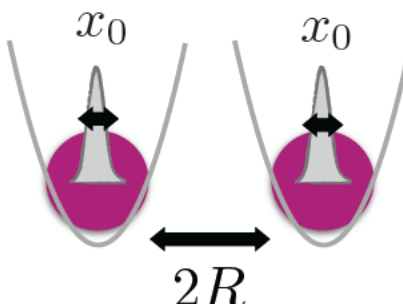
Chapel Hill Conference 1957 (29)

WITTEN: “What prevents this from becoming a practical experiment?”

An ultimate experiment? Entanglement by gravity...



Entangling via gravity



Entangling time:

$$t_{ent} = \frac{2\hbar}{GM^2} R^3 \left(\frac{1}{x_0}\right)^2$$

Pino et al., arXiv:1603.01553

$|12\rangle_{down}$

FEYNMAN: "Therefore, provided that the duce a gravitation destroy the possi is a bare possibilit fails and becomes because of some chain. But aside

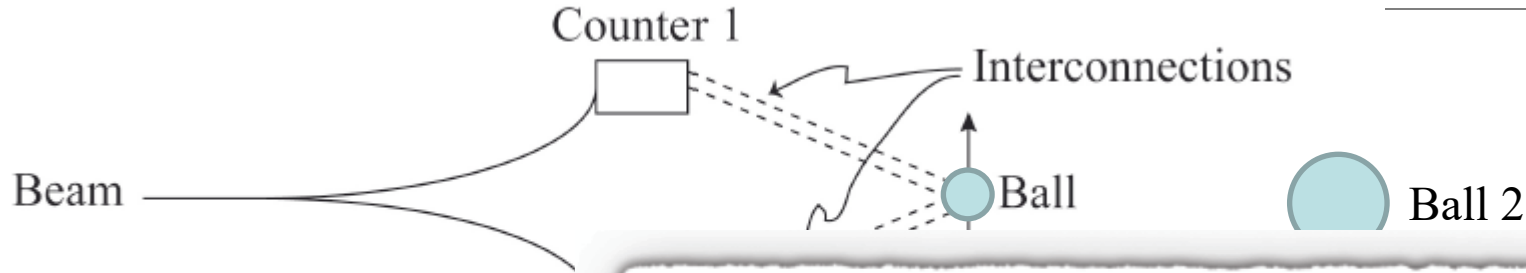
ics up to any level then you have to believe in gravitational quantization in order to describe this experiment."

Chapel Hill Conference 1957 (29)

Example: For 2 **lead spheres** of **diameter 500 μm**, an initial **superposition size for sphere 1** of $\Delta r = 5 \times 10^{-7}$ m and preparation of **sphere 2** in a **motional ground state** (100 Hz trap frequency) with $\Delta x_0 = 10^{-15}$ m, we obtain $\Gamma_{ent} = 1.5$ Hz, i.e. **gravitational entanglement** is established on a **second time scale**.

$$\Gamma_{ent} = \left(\frac{GM}{\hbar}\right) \Delta r \rho \Delta x_0$$

An ultimate experiment? Entanglement by gravity...



Entangling via gravity

Entangling time:

$$t_{ent} = \frac{2\hbar}{GM^2} R^3 \left(\frac{1}{x_0}\right)^2$$

Pino et al., arXiv:1603.01553

$|12\rangle_{down}$

FEYNMAN: "Therefore, provided that the duce a gravitation destroy the possi is a bare possibilit fails and becomes because of some chain. But aside

ics up to any level then you have to believe in gravitational quantization in order to describe this

Note: **dynamical potential landscape** allows for significantly smaller masses (Pino 2016: **2 μm**)

Example: For 2 **lead spheres** of diam **superposition size for sphere 1** of $\Delta r = 5 \times 10^{-7}$ m and **of sphere 2 in a motional ground state** (100 Hz trap frequency) with $\Delta x_0 = 10^{-15}$ m, we obtain $\Gamma_{ent} = 1.5$ Hz, i.e. **gravitational entanglement is established on a second time scale.**

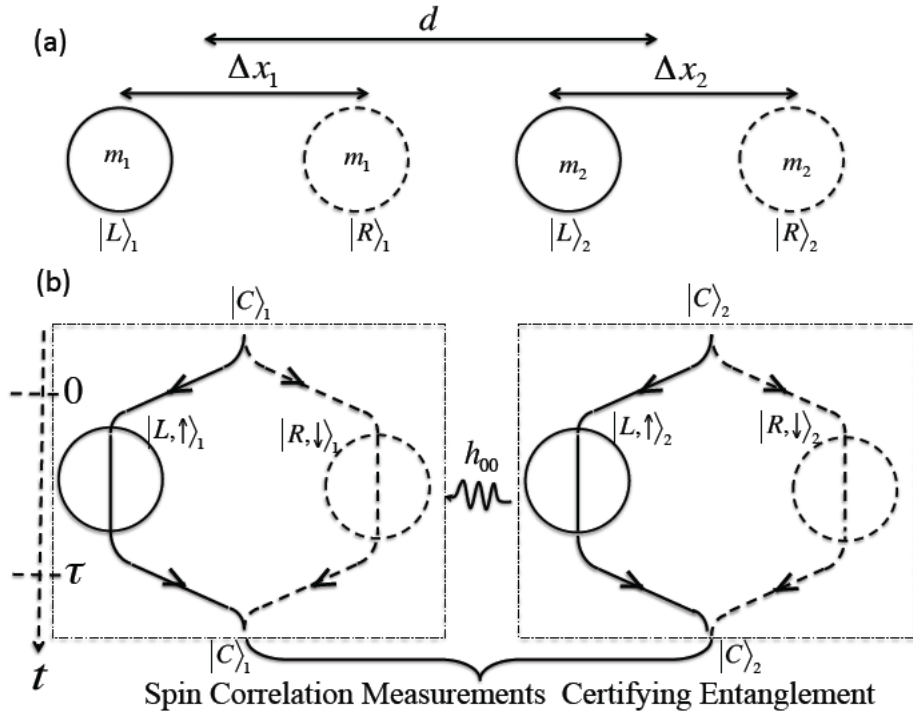
Hill Conference 1957 (29)

$$\Gamma_{ent} = \left(\frac{GM}{\hbar}\right) \Delta r \rho \Delta x_0$$

An ultimate experiment? Entanglement by gravity...

Refined proposal by *Bose, Kim, Milburn et al. 2017*:
Entanglement by gravitational phase shift (COW) and CSIGN gate

Beam 



FEYNMAN: "Therefore provided that produce a gravita destroy the p is a bare possi fails and beco because of so chain. But as

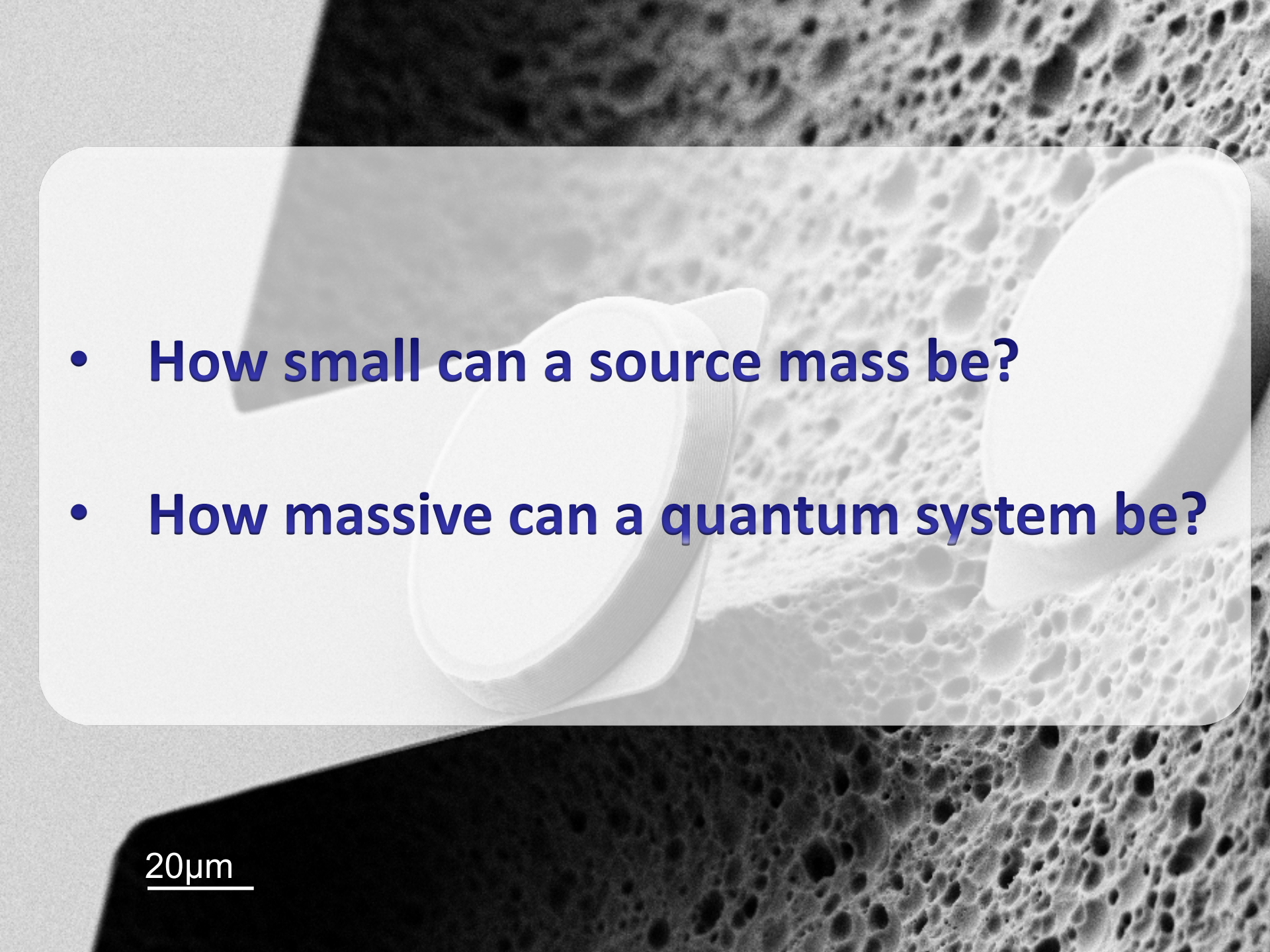
Bose et al., PRL 119, 240401 (2017)

Note: **dynamical potential landscape** allows for significantly smaller masses (*Pino 2016: 2 μm*)

Hill Conference 1957 (29)

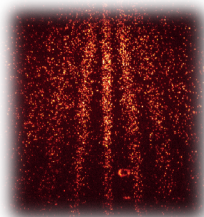
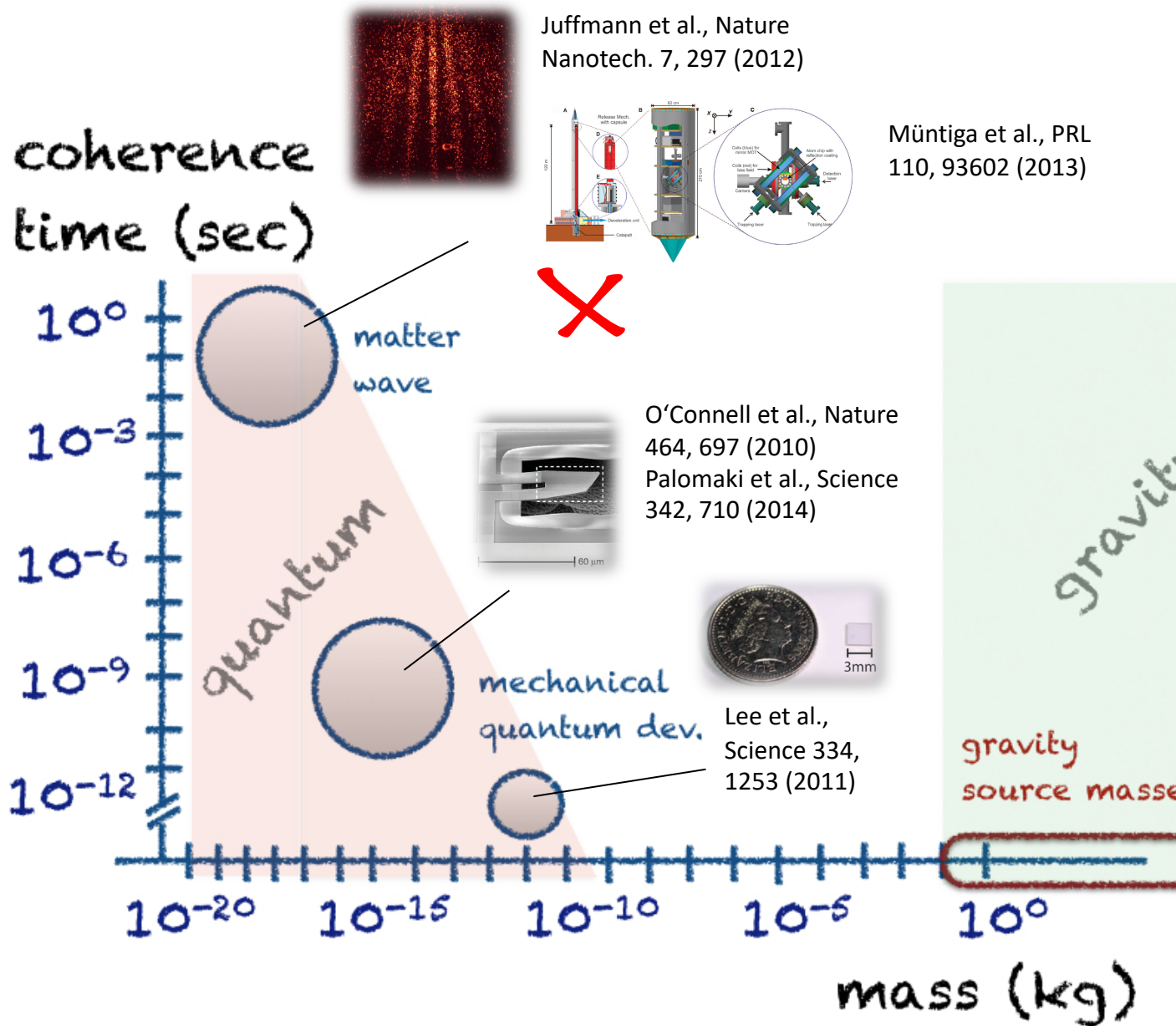
Example: For 2 lead spheres of diam superposition size for sphere 1 of $\Delta r = 5 \times 10^{-7}$ m and of sphere 2 in a motional ground state (100 Hz trap frequency) with $\Delta x_0 = 10^{-15}$ m, we obtain $\Gamma_{\text{ent}} = 1.5$ Hz, i.e. gravitational entanglement is established on a second time scale.

$$\Gamma_{\text{ent}} = \left(\frac{GM}{\hbar}\right) \Delta r \rho \Delta x_0$$

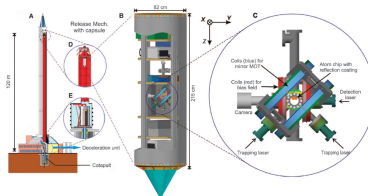
- 
- A scanning electron micrograph (SEM) showing two white, pill-shaped structures on a porous, textured surface. The structures are oval-shaped with a slightly raised rim. The background is a complex, porous network of interconnected fibers or cells. A scale bar in the bottom left corner indicates 20 micrometers.
- **How small can a source mass be?**
 - **How massive can a quantum system be?**

20 μ m

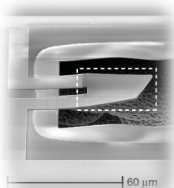
How massive/small can we go?



Juffmann et al., Nature Nanotech. 7, 297 (2012)



Müntiga et al., PRL 110, 93602 (2013)



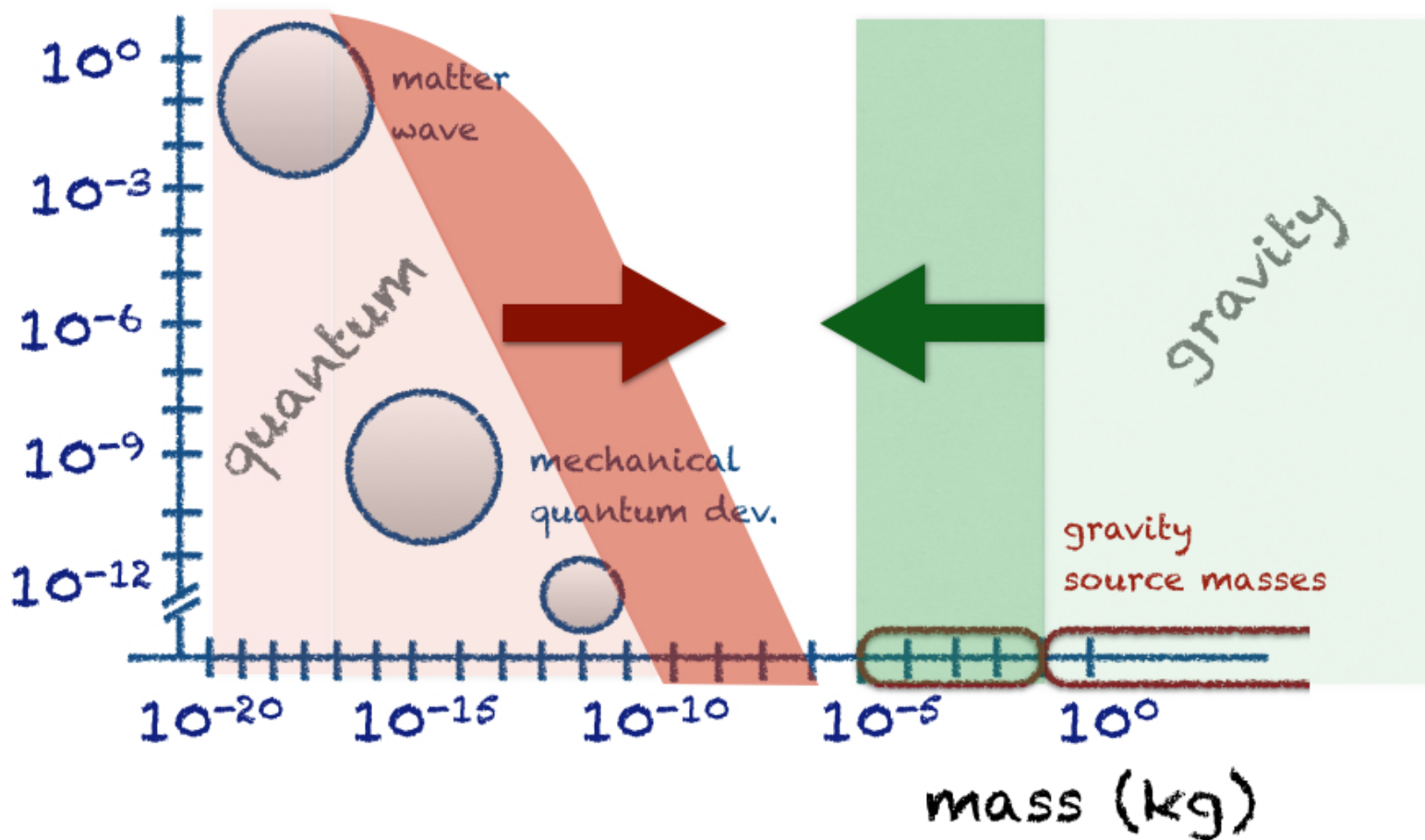
O'Connell et al., Nature 464, 697 (2010)
Palomaki et al., Science 342, 710 (2014)



Lee et al., Science 334, 1253 (2011)

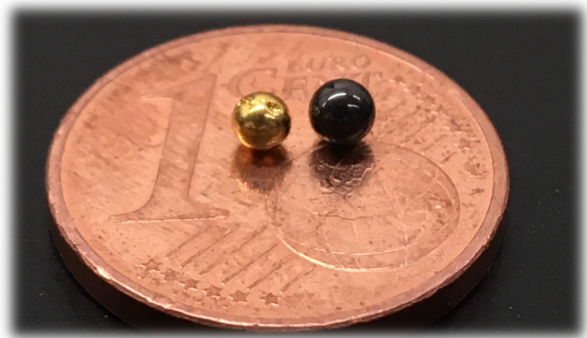
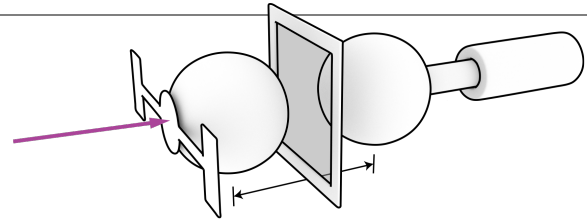
How massive/small can we go?

coherence
time (sec)



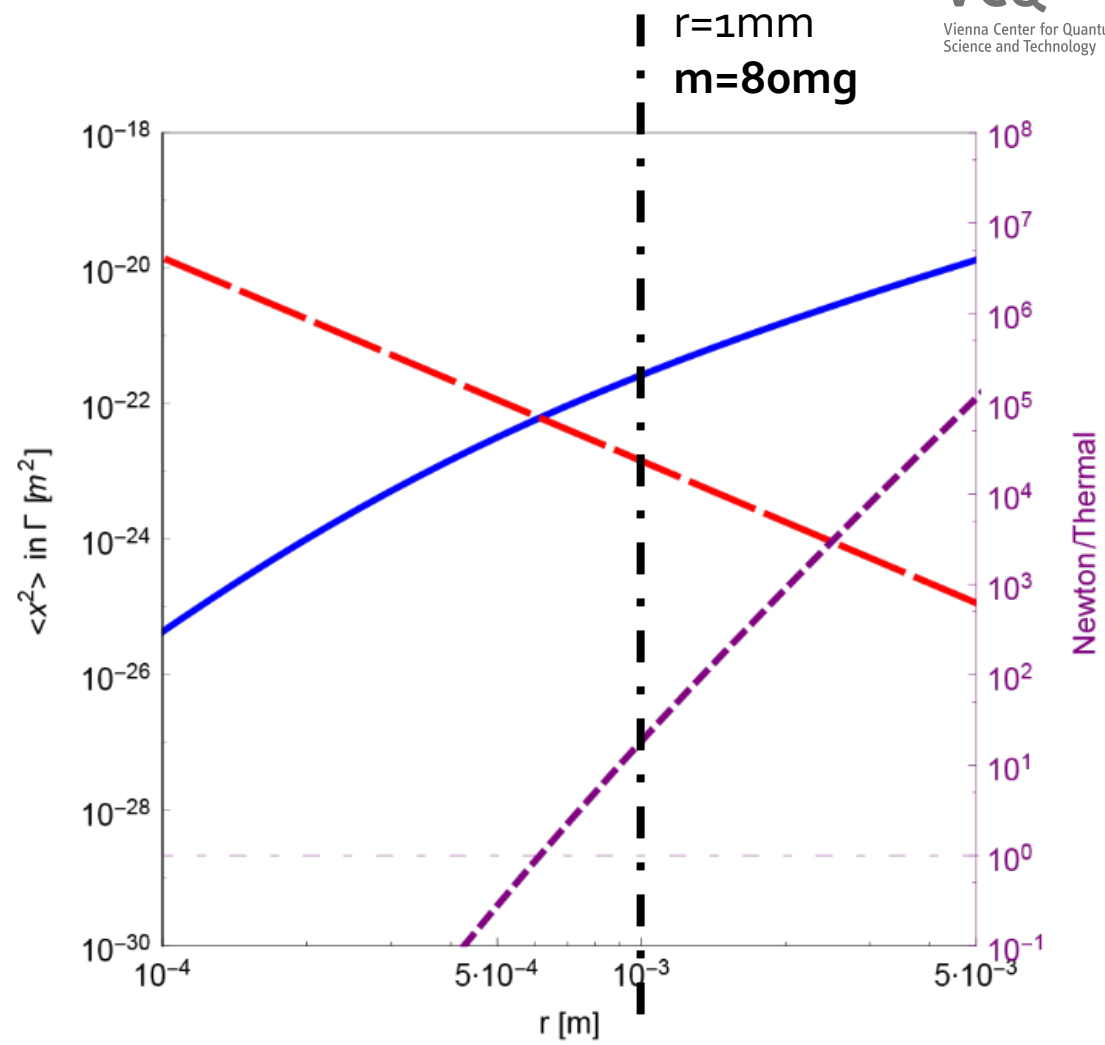
Measuring gravity between microscopic source masses ?

Schmöle et al., *Class. Quant. Grav.* 33, 125031 (2016)



In collaboration with Gröblacher group (TU Delft)

© Scientific American



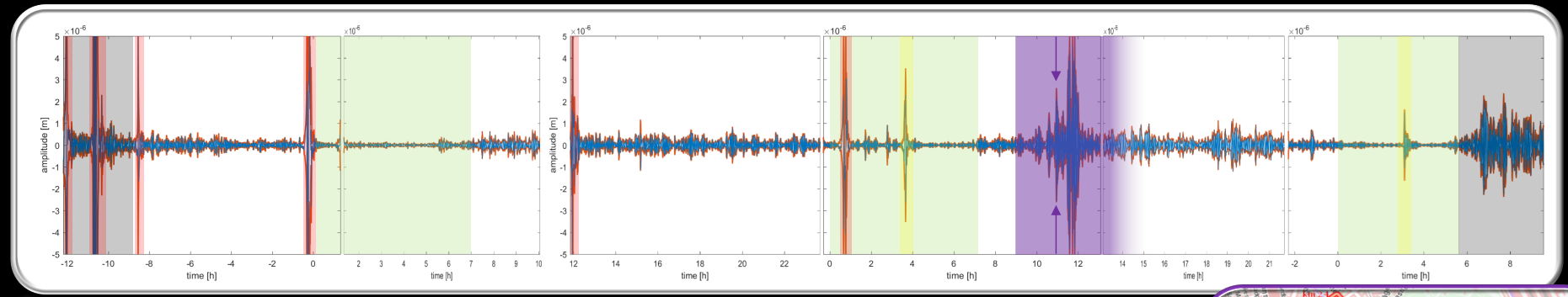
Smallest source mass to date: **0.7 g**

Mitrofanov & Ponomareva, *Zh. Eksp. Teor. Fiz.* 94,16-22 (1988)

The vibration isolation challenge...



1 weekend in the life of milli-g



↑ Friday 5.4. 12:00 ↑ 00:00 ↑ Saturday 6.4. 12:00 ↑ 00:00 ↑ Sunday 7.4. 12:00

- 1 weekend time-series of milli-tor
- Discharging experiment
- Status: damped & DC-locked
- 3mHz-100mHz band
- Signal & instantaneous amplitude

- Table or pendulum work
- Nighttime (0:00-7:00/0:30-5:30)
- Normal weekday
- Something new
- Marathon (1st finisher)



Big G: the open problem

The search for

Newton's constant

Clive Speake and Terry Quinn



The "G machine," now housed at the University of Birmingham in the UK, was used at the International Bureau of Weights and Measures in France to measure Newton's gravitational constant.

Three decades of careful experimentation have painted a surprisingly hazy picture of the constant governing the most familiar force on Earth.

Physics Today July 2014

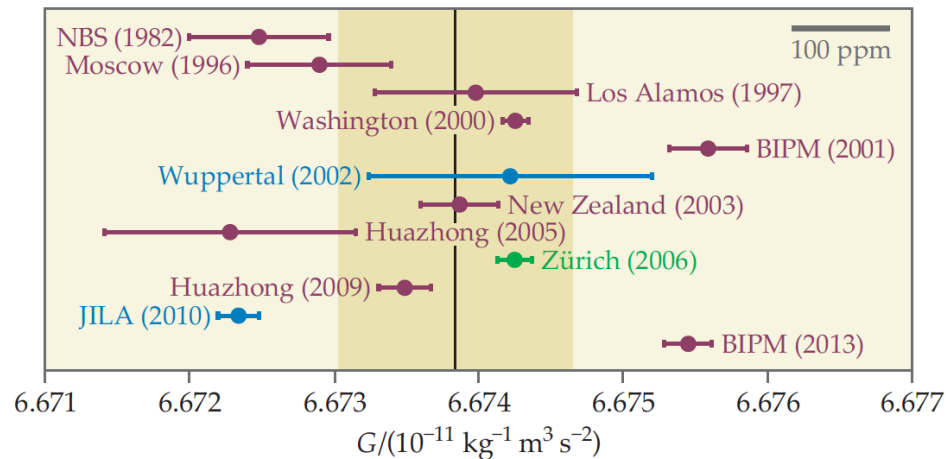


Figure 1. Measurements of Newton's gravitational constant G have yielded conflicting results. Here, the results of torsion-balance (maroon), pendulum (blue), and beam-balance (green) experiments discussed in the text are shown, along with the location and year in which they were measured. Error bars correspond to one standard deviation; the shaded region indicates the assigned uncertainty of the value recommended by the Committee on Data for Science and Technology in 2010. (Adapted from T. J. Quinn et al., *Phys. Rev. Lett.* **111**, 101102, 2013.)

NEWS

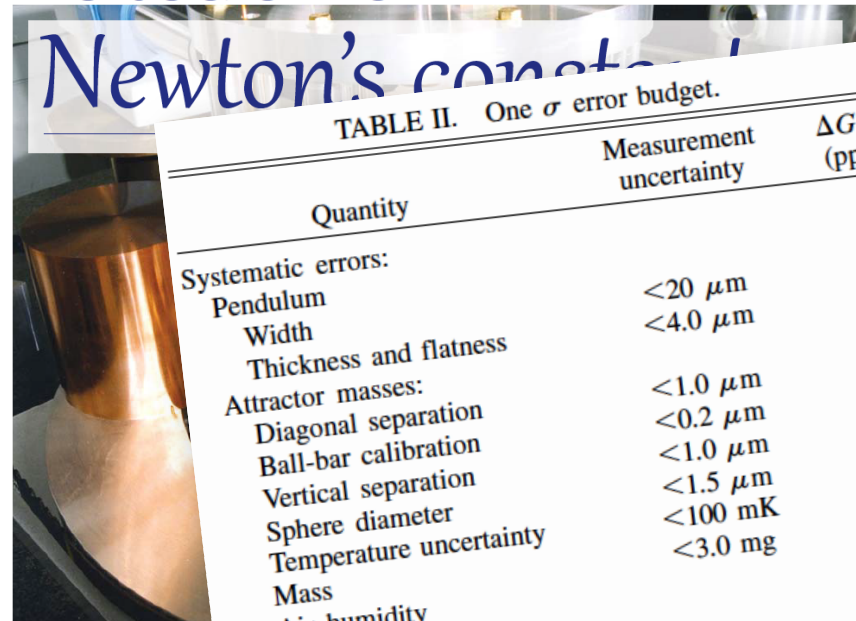
NATURE | Vol 466 | 26 August 2010

G-whizzes disagree over gravity

Recent measurements of gravitational constant increase uncertainty over accepted value.

Big G: the open problem

The search for



Three decades of carefully taken, but very hazy picture of the copper pendulum

Newton's constant

TABLE II. One σ error budget.

Quantity	Measurement uncertainty	$\Delta G/G$ (ppm)
Systematic errors:		0.4
Pendulum	$<20 \mu\text{m}$	4.0
Width	$<4.0 \mu\text{m}$	
Thickness and flatness		7.1
Attractor masses:	$<1.0 \mu\text{m}$	1.4
Diagonal separation	$<0.2 \mu\text{m}$	5.2
Ball-bar calibration	$<1.0 \mu\text{m}$	2.6
Vertical separation	$<1.5 \mu\text{m}$	6.9
Sphere diameter	$<100 \text{mK}$	0.4
Temperature uncertainty	$<3.0 \text{mg}$	0.5
Mass		0.3
Air humidity		0.6
Residual twist angle		0.4
Magnetic fields		0.1
Rotating temperature gradient	$<10^{-7}$	2.0
Time base		5.8
Data reduction		13.7
Statistical error:		
Total:		

Physics Today July 2014

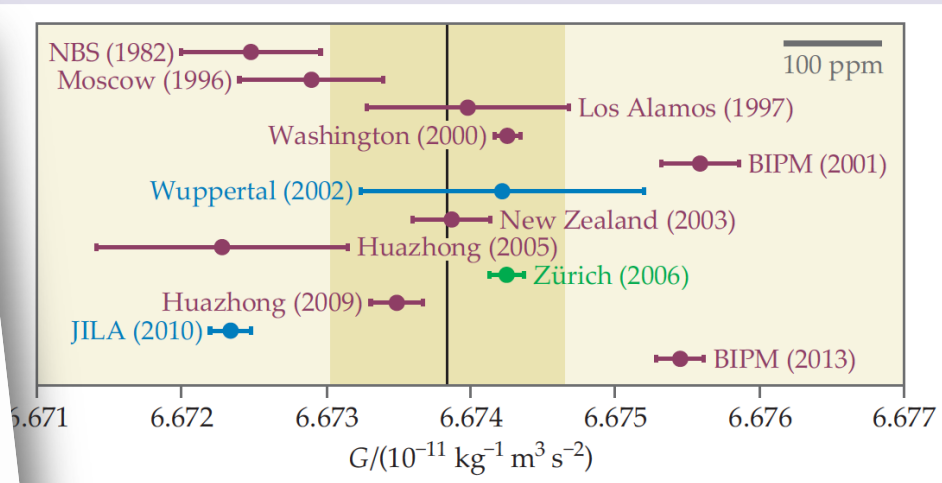


Figure 1. Measurements of Newton's gravitational constant G have yielded conflicting results. Here, the results of torsion-balance (maroon), pendulum (blue), and beam-balance (green) experiments discussed in the text are shown, along with the location and year in which they were measured. Error bars correspond to one standard deviation; the shaded region indicates the assigned uncertainty of the value recommended by the Committee on Data for Science and Technology in 2010. (Adapted from T. J. Quinn et al., *Phys. Rev. Lett.* **111**, 101102, 2013.)

NEWS

G-whizzes disagree over gravity

Recent measurements of gravitational constant increase uncertainty over accepted value.

NATURE | Vol 466 | 26 August 2010

Big G: the open problem

Physics Today July 2014

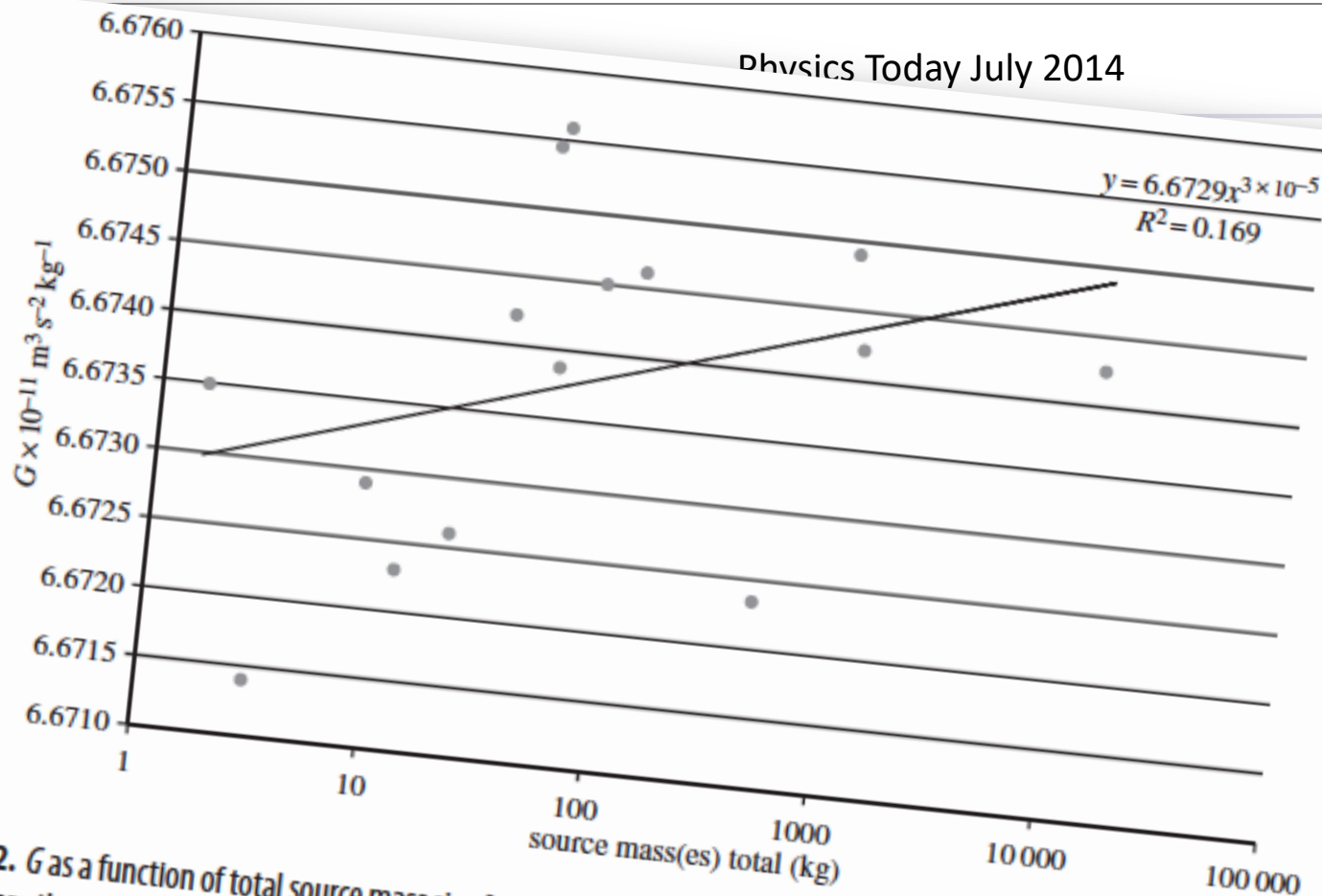


Figure 2. G as a function of total source mass size for the measurements with $\Delta G/G < 250$ ppm. The 15 data points from left to right are the results from Tu *et al.* [12], Pontikis [13], Karagioz *et al.* [15], Hu *et al.* [18], Luther *et al.* [23], Gundlach *et al.* [25], Quinn *et al.* [27], Quinn *et al.* [28], Armstrong *et al.* [29], Sagitov *et al.* [30], R. D. Newman (2013, personal communication), Parks *et al.* [37], Nolting *et al.* [44], Kleinvoß [45] and Schlamminger *et al.* [47].

From: G. T. Gillies, C. S. Unnikrishnan, *Phil. Trans. R. Soc. A* 372:20140022 (2014)

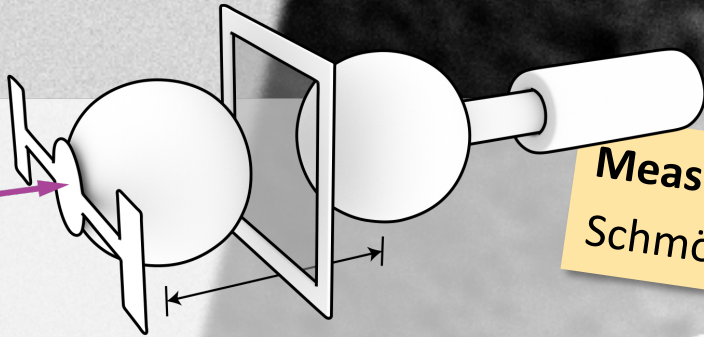


constant G have
-balance (maroon),
ments discussed in
in which they were
eviation; the shaded
ue recommended
ogy in 2010.
(101102, 2013.)

Thru
hoi

The
N

...ments of gravitational constant increase uncertainty over accepted value.

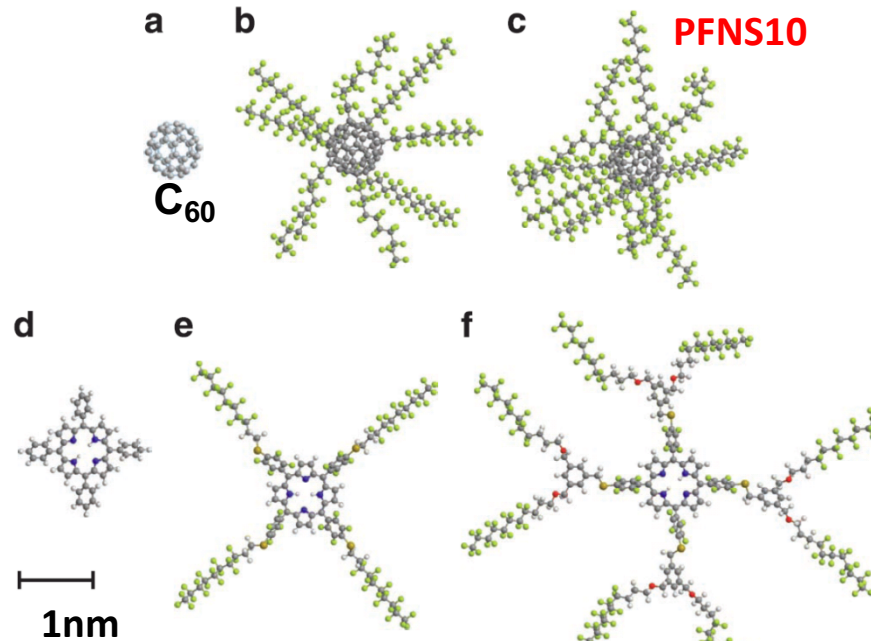
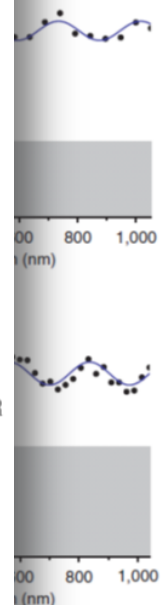
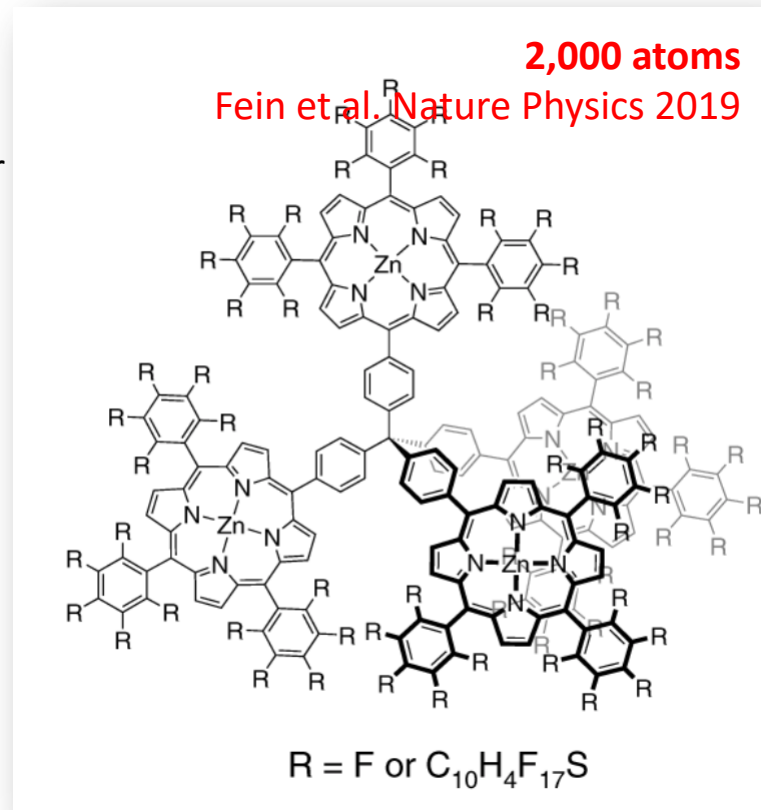
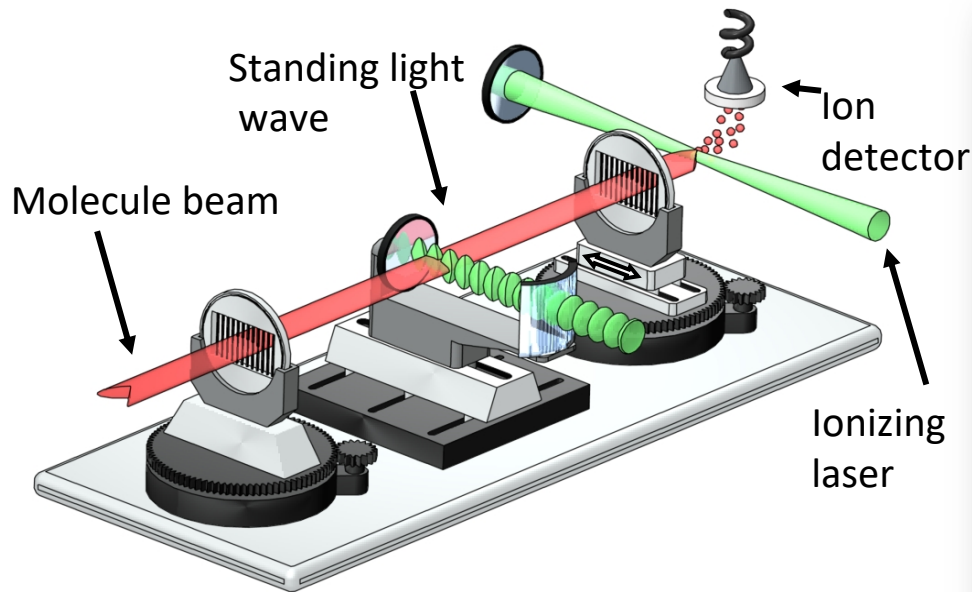


Measuring the gravitational force of milligram masses
Schmöle et al., *Class. Quant. Grav.* **33**, 125031 (2016)

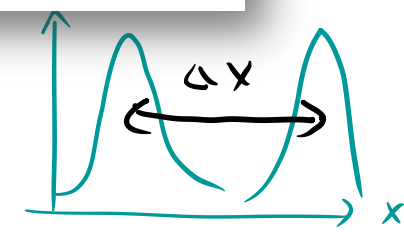
- **How small can a source mass be?**
- **How massive can a quantum system be?**

20 μ m

COM superposition states of massive systems: where do we stand?

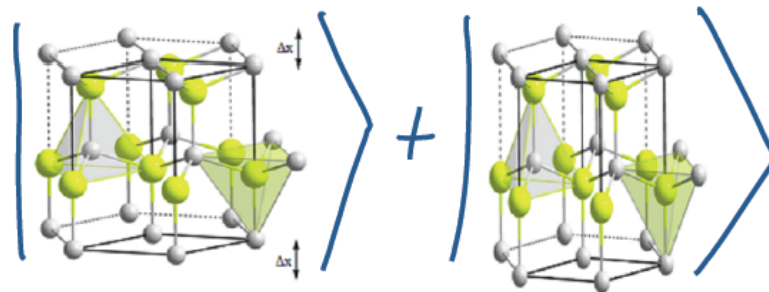


PFNS10: C₆₀[C₁₂F₂₅]₁₀
(perfluoroalkylated nanosphere)
430 atoms
 $m \sim 10^{-23} \text{ kg} = 6910 \text{ AMU}$
 $\Delta x \sim 100 \text{ nm}$ (~50x its diameter)



COM superposition states of massive systems: where do we stand?

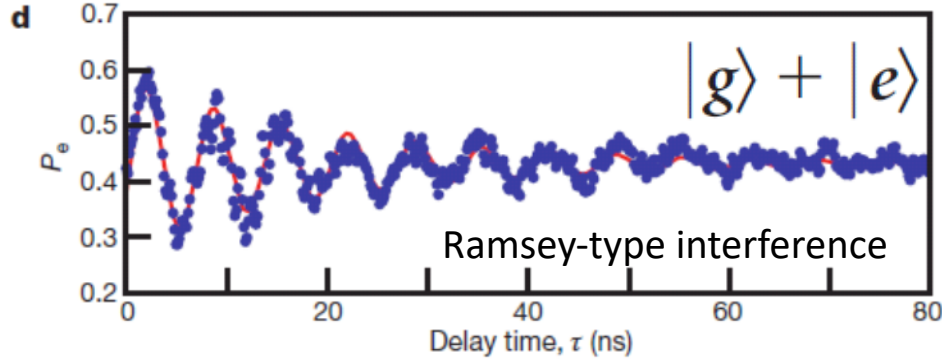
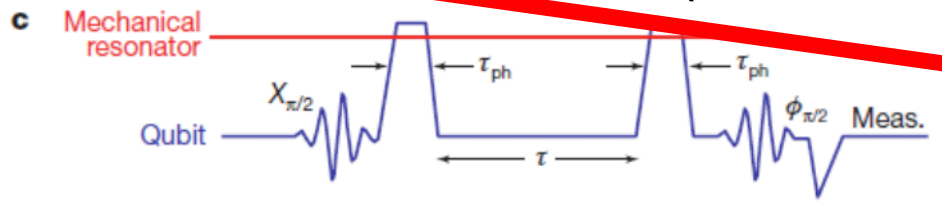
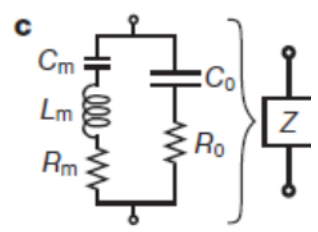
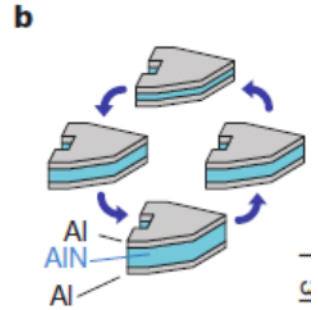
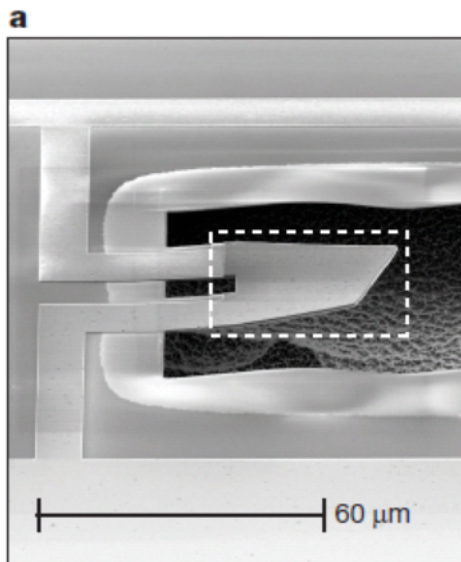
Micromechanics, 2×10^{13} atoms
 $m \sim 10^{-12} \text{ kg} = 7 \times 10^{14} \text{ AMU}$
 $\Delta x \sim 10^{-16} \text{ m}$ ($\sim 10^{-10}$ x its diameter)



Quantum ground state and single phonon control of a mechanical resonator

A. D. O'Connell¹, M. Hofheinz¹, M. Ansmann¹, Radoslaw C. Bialczak¹, M. Lenander¹,
 D. Sank¹, H. Wang¹, M. Weides¹, J. Wenner¹, John M. Martinis¹ & A. N. Cleland¹

Observed coherence time ($>20\text{ns}$)
 bounds **Penrose's "mass distribution localization"** to $R > 10^{-21} \text{m}$



6 GHz thickness oscillation
 $\rightarrow n \sim 0.07$ @ 20 mK

Note: $E_g - E_e = h \cdot f_m \approx 20 \mu\text{eV}$

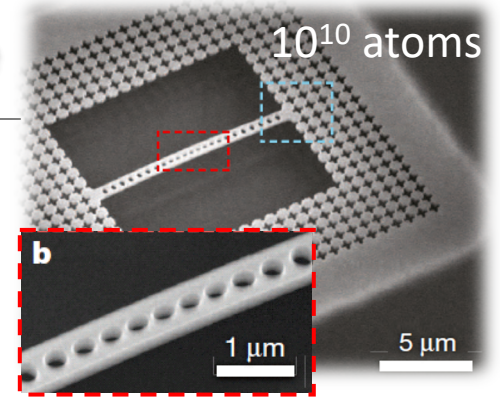
Mechanical Systems IN the quantum regime

Quantum ground state of motion

Microwave cavity cooling: Teufel et al., Nature 475, 359 (2011)

Laser cooling: Chan et al., Nature 478, 89 (2011)

... and many more around the world...



Quantum squeezed states of motion

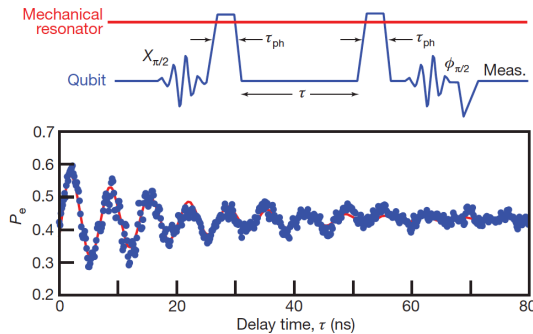
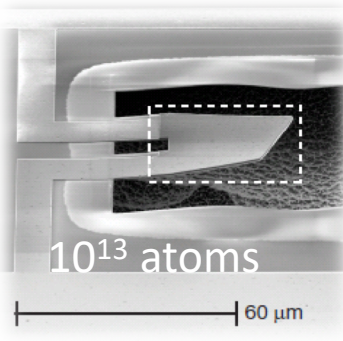
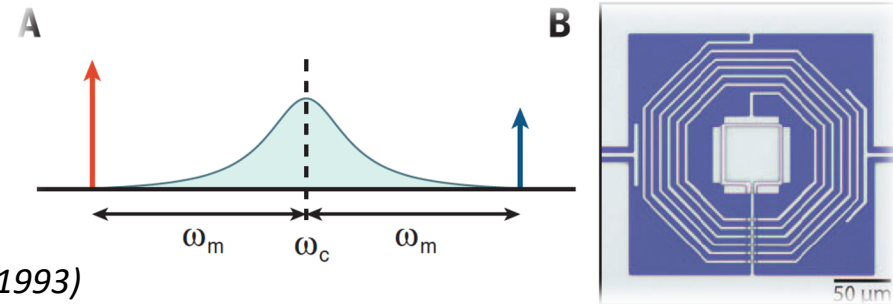
Wollman et al., Science 349, 952 (2015)

J.-M. Pirkkalainen et al., PRL 115, 243601 (2015)

F. Lecocq et al., PRX 5, 041037 (2015)

„reservoir engineering“

(see also Cirac et al. PRL 70, 556 1993)



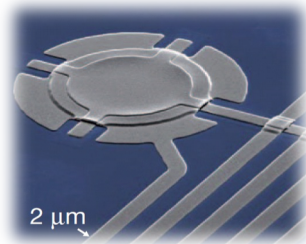
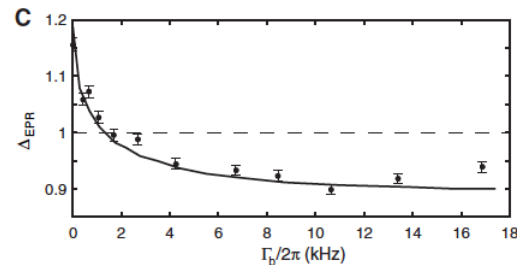
Non-Gaussian quantum states of motion

Phonon control through superconducting qubit:

O’Connell et al., Nature 464, 697 (2010)

Photon-phonon correlations:

Riedinger, Hong et al., Nature 530, 313 (2016)



Quantum entanglement

EPR-type entanglement (MW):

Palomaki et al., Science 342, 710 (2013)

Bell-type entanglement (optical):

Lee et al., Science 334, 1253 (2011)

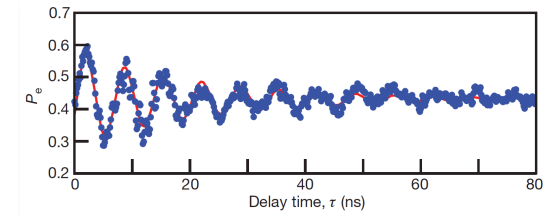
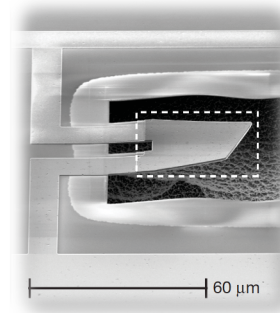
Riedinger et al., Nature 556, 473 (2018)

Q: How to achieve large mass **AND** long coherence time in a quantum experiment?

Solid-state mechanical quantum devices
(clamped):

$10^{10} - 10^{16}$ atoms

Coherence time τ_c $10^{-12} - 10^{-8}$ sec

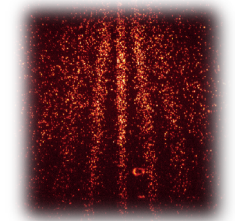
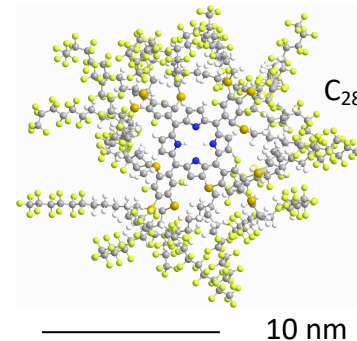


O'Connell et al., Nature 464, 697 (2010)

Matter-wave interferometry (free-fall):

$10^0 - 10^4$ atoms

Coherence time τ_c $10^{-3} - 10^0$ sec



Juffmann et al., Nature
Nanotech. 7, 297 (2012)

A: Quantum control of levitated mechanical systems!



- Quantum control of a trapped massive object $\gg 10^{10}$ atoms
- Long coherence times (up to seconds)
- Exceptional force sensitivity
- **Externally engineerable (and controllable) arbitrary potential landscape**

Optically levitating nanoparticles

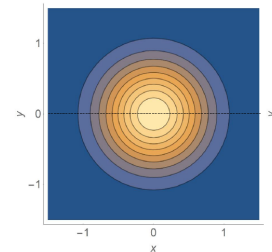
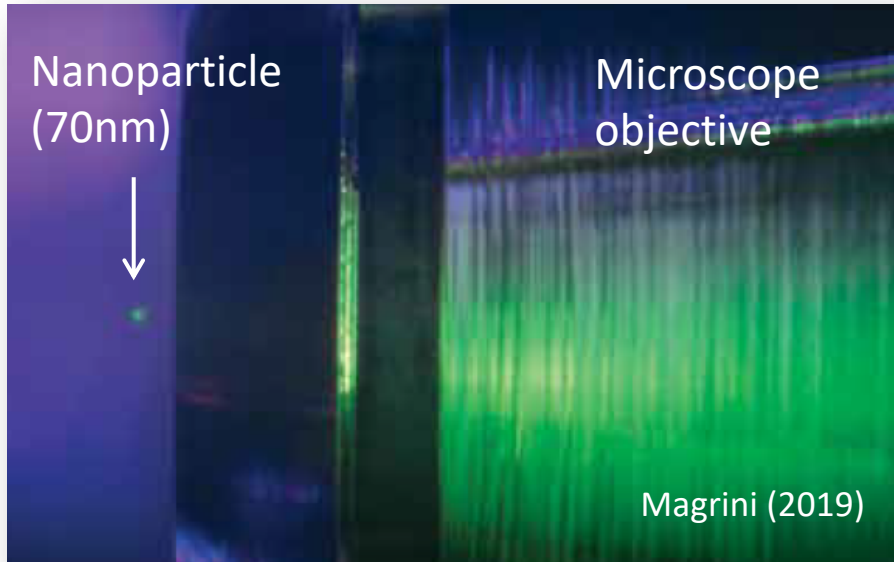
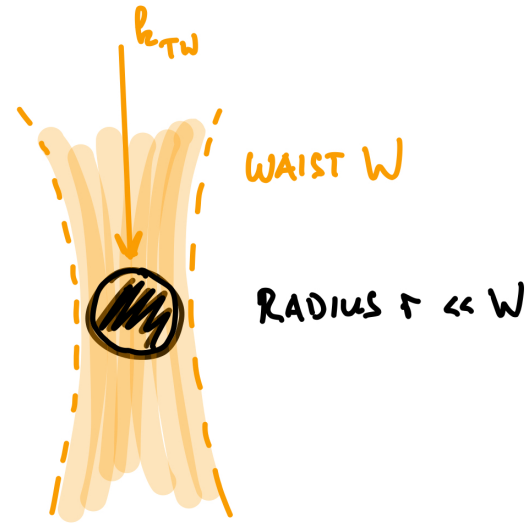
OPTICAL LEVITATION:

$$\hat{H} < \underline{d} \cdot \underline{E} = \alpha \cdot E^2$$

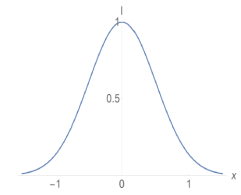
α : $\text{Re}\{\text{Polarizability}\}$
 E : optical trapping field

\hookrightarrow beam intensity

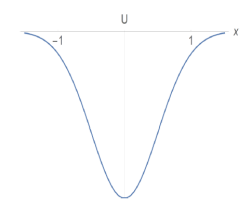
$$\rightarrow \text{GRADIENT FORCE } F \propto (\nabla E^2) \cdot \underline{\alpha}$$



$\alpha = 0$



intensity

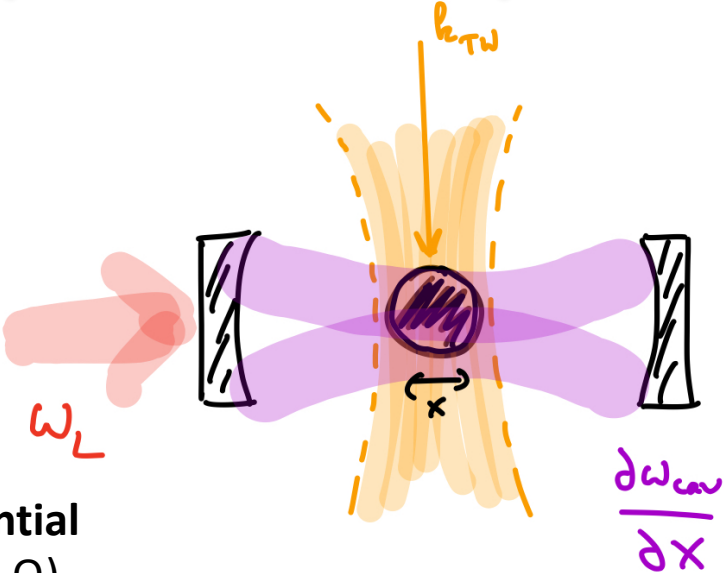


potential

Pioneering work by Ashkin:
 A. Ashkin, PRL 24, 156 (1970).
 A. Ashkin, J. M. Dziedzic, APL 28, 333 (1976).

Towards quantum state preparation of a free particle

Chang et al., PNAS 2010
 Romero-Isart et al., NJP 2010
 P. F. Barker et al., PRA 2010
early work:
 Hechenblaikner, Ritsch et al., PRA 58, 3030 (1998)
 Vuletic & Chu, PRL 84, 3787 (2000)

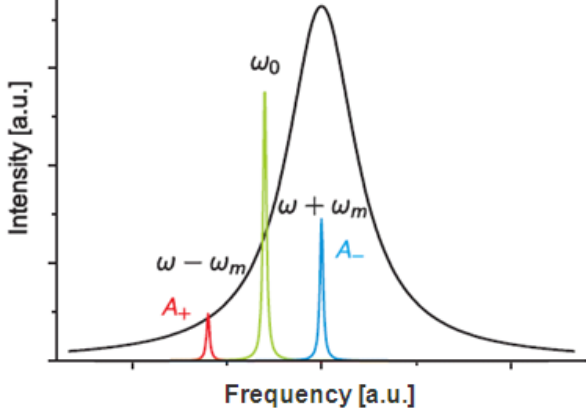


- Harmonic oscillator in optical potential (negligible support loss, high Q)
- Quantum control via cavity optomechanics (laser cooling, state transfer, etc.)

Cavity-Optomechanics

Rev.Mod.Phys. 86, 1391 (2014)

pioneering works by V. Braginsky



OPTOMECHANICS:

$$\hat{H}_{int} \propto \alpha \cdot E_{cav}^2 = \hbar g_0 \underbrace{\hat{a}^\dagger \hat{a}}_{\text{cavity}} \underbrace{(\hat{d} + \hat{d}^\dagger)}_{\text{mechanics}}$$


with $g_0 = \frac{\partial \omega_c}{\partial x} = \left[\frac{\alpha}{V_c} \right] \frac{\omega_c^2}{c \epsilon_0} x_{zfp}$

α : Re { polarizability }
 V_c : cavity mode volume
 x_{zfp} : zero-point motion


Cavity control of levitated particles

- dielectrics and superconductors (100nm – 10 μ m)
- ultimate goal: anharmonic coupling (e.g. via 2-level systems) & full quantum control

c *photonic crystal cavity*

2 μ m 

optical fiber

- 
- Grass et al., Appl. Phys. Lett. 108, 221103 (2016): optical control in hollow-core fibres
Kiesel et al., PNAS 110, 14180 (2013): cavity cooling of levitated particles
Magrini et al., Optica 5, 1597 (2018): near-field coupling to photonic crystal cavity
Delic et al., PRL 122, 123602 (2019): cavity cooling via coherent scattering

Cavity Optomechanics with levitated nanoparticles

Ashkin since 1967

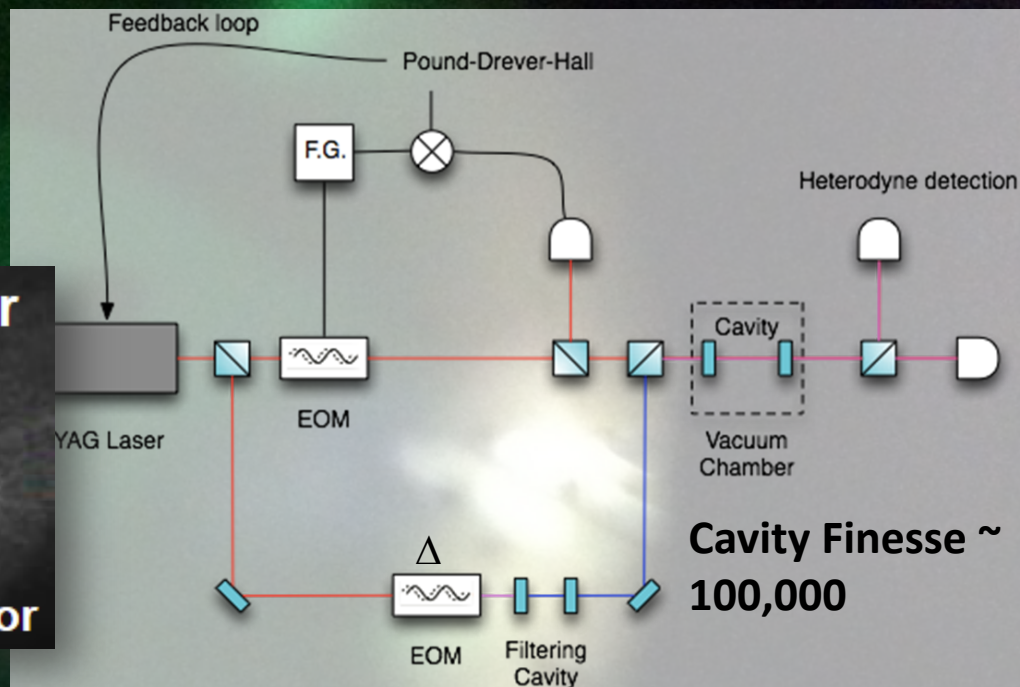
Raizen group, *Science* 2010

Novotny, Quidan 2012

Barker group 2014

Geraci group 2015

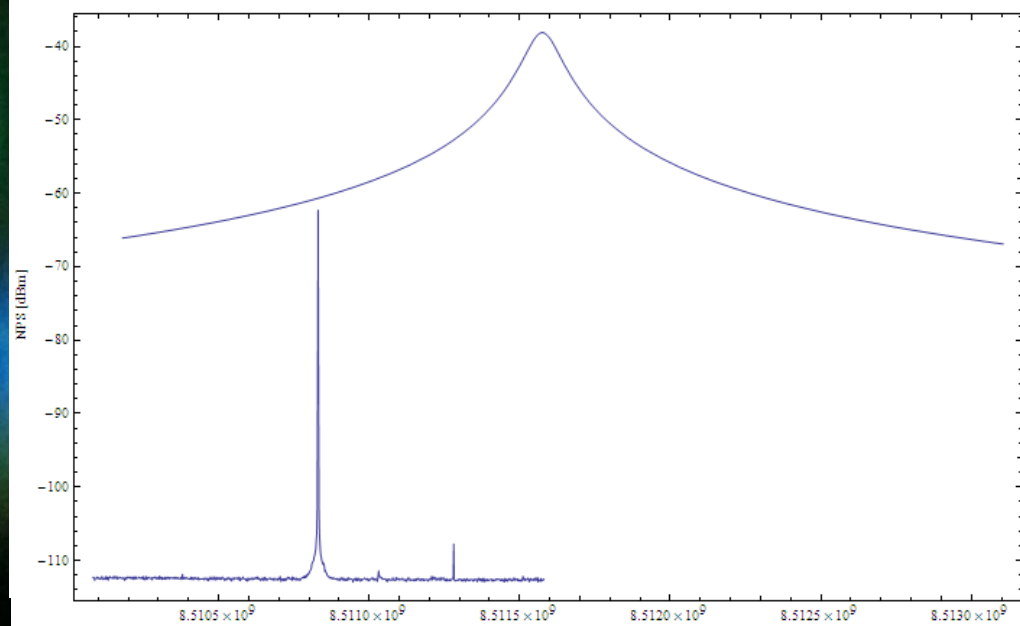
Levitation in Cavity @ 4mbar



Optical trapping and cavity cooling ($R \sim 200\text{nm}$)

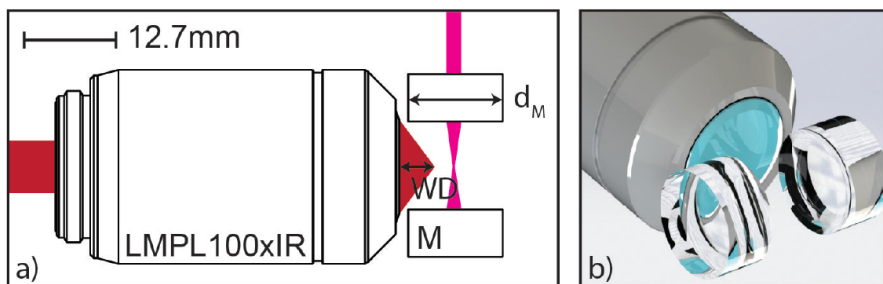
Kiesel, Delic, Grass et al.,

PNAS USA 110, 14180 (2013)



Cavity Optomechanics with levitated nanoparticles

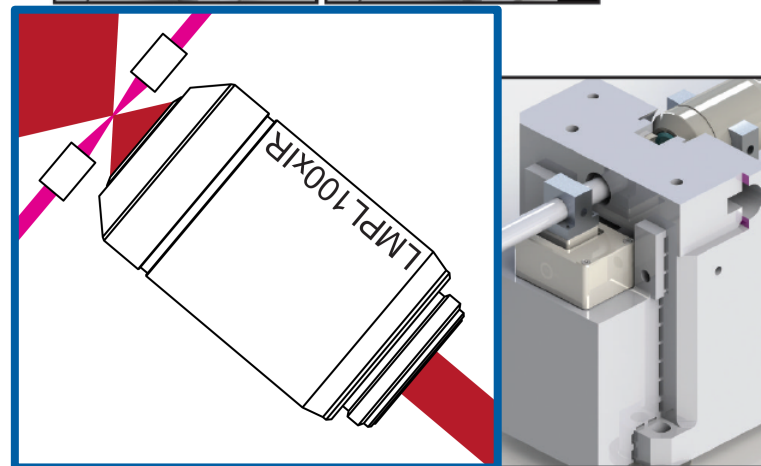
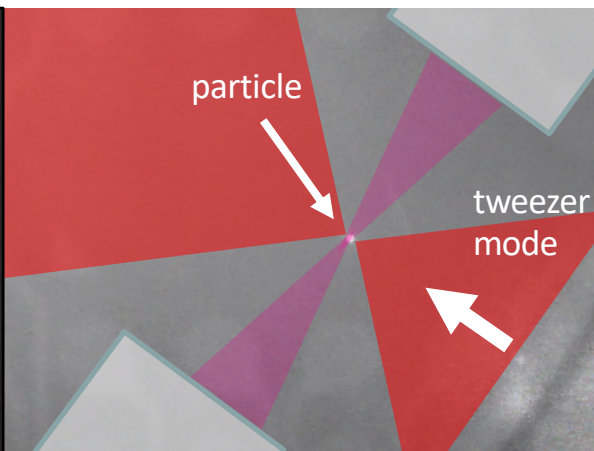
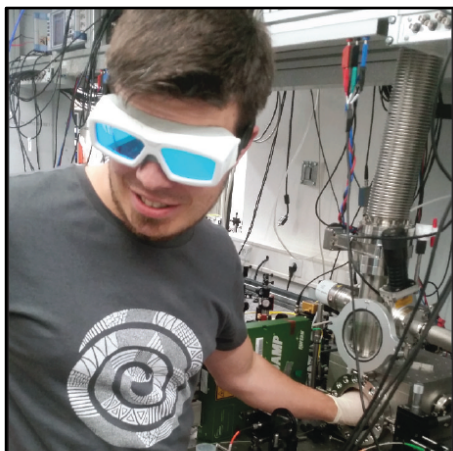
Delić, Grass et al., arXiv:1902.06605



Working distance of tweezer < Radius of cavity mirrors
→ mirror cutting (@Weitz group, U Bonn)

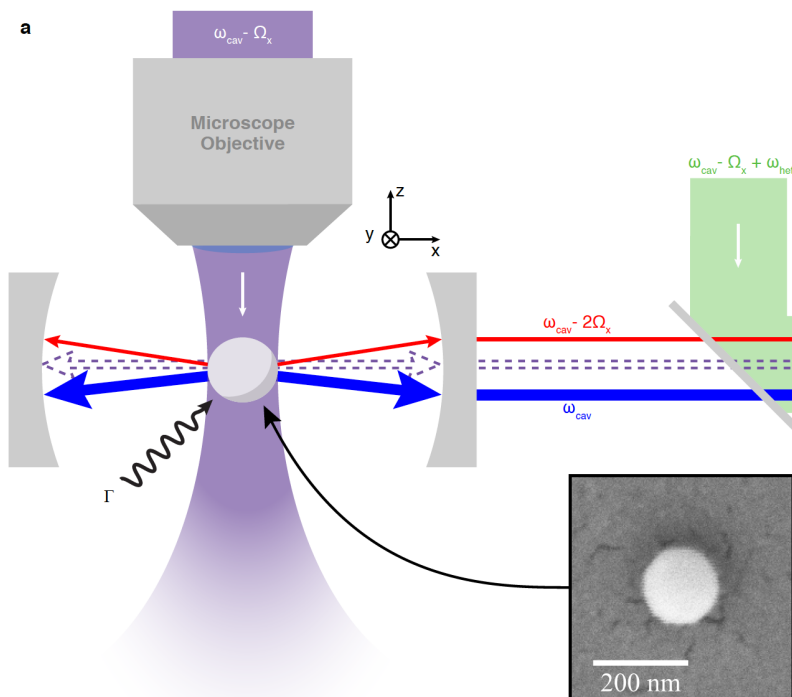


a) b)

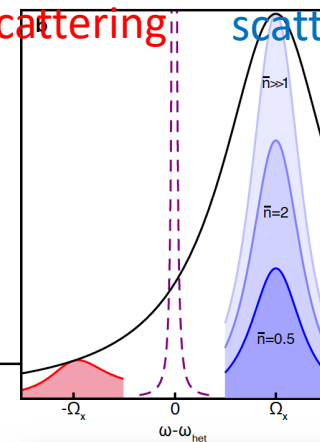


Cavity Optomechanics with levitated nanoparticles

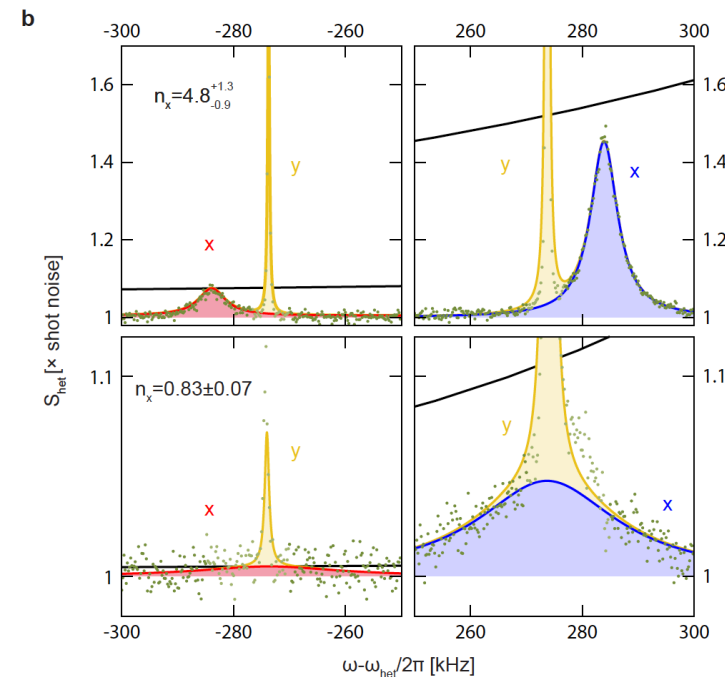
Delić et al., arXiv:1911.04406



Stokes scattering
anti-Stokes scattering

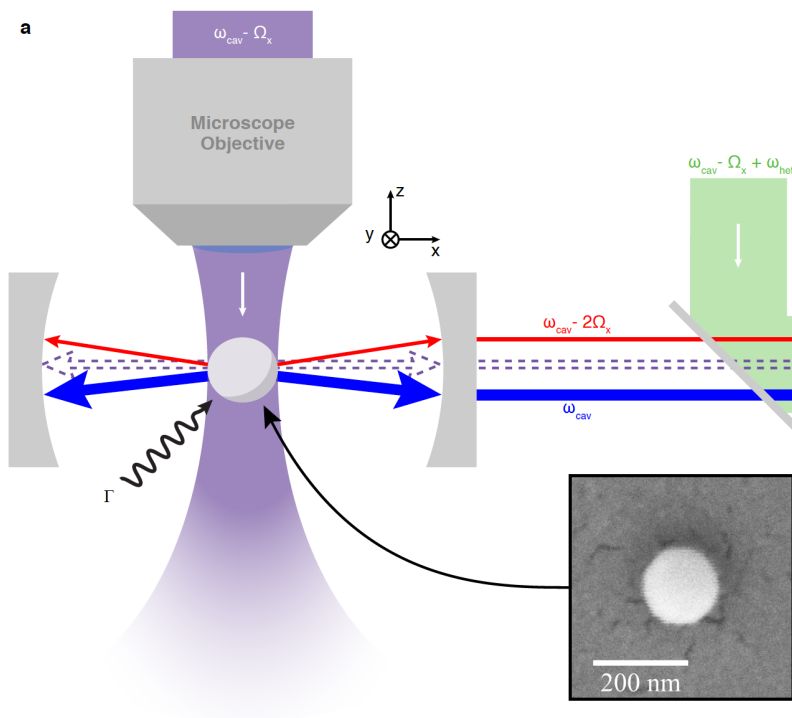


$$\begin{aligned} \omega_x &\approx 2\pi \times 305 \text{ kHz} \\ \omega_z &\approx 2\pi \times 80 \text{ kHz} \\ \omega_y &\approx 2\pi \times 275 \text{ kHz} \\ \kappa &= 2\pi \times 193 \text{ kHz} \\ p &= 1e-6 \text{ mbar}, T = 300\text{K} \end{aligned}$$

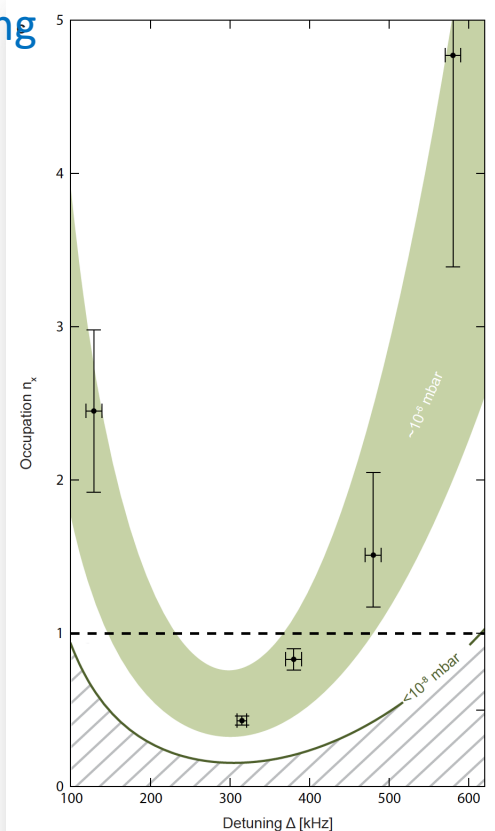
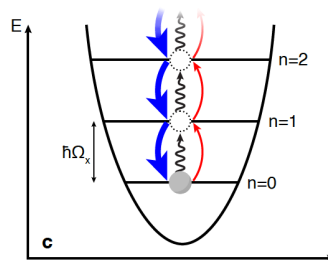
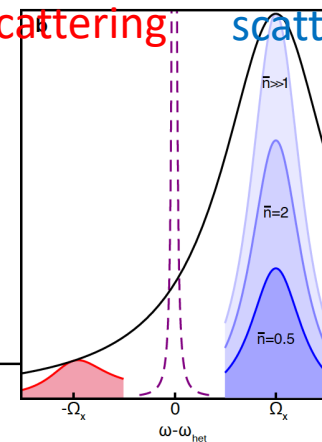


Cavity Optomechanics with levitated nanoparticles

Delić et al., arXiv:1911.04406



Stokes scattering anti-Stokes scattering



$$\longleftrightarrow \omega_x \approx 2\pi \times 305 \text{ kHz}$$

$$\updownarrow \omega_z \approx 2\pi \times 80 \text{ kHz}$$

$$\swarrow \omega_y \approx 2\pi \times 275 \text{ kHz}$$

$$\kappa = 2\pi \times 193 \text{ kHz}$$

$$p = 1e-6 \text{ mbar}, T = 300\text{K}$$

$n_x < 0.5$ (ground state probability $> 2/3$)
 Center-of-mass $T_c = 12\mu\text{K}$; environment $T_e > 300\text{K}$
 $g_x = 2\pi \times 71 \text{ kHz}$, Cooperativity $C = 5$

Cavity Optomechanics with levitated nanoparticles

Delić et al., arXiv:1911.04406

Limitations:

- Sideband resolution
- Gas and recoil scattering
- Phase noise

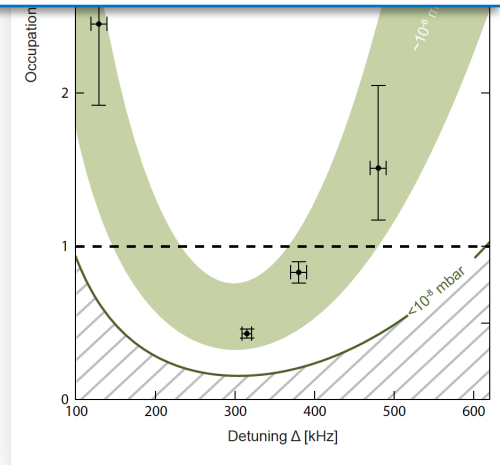
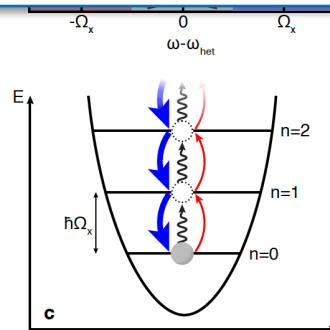
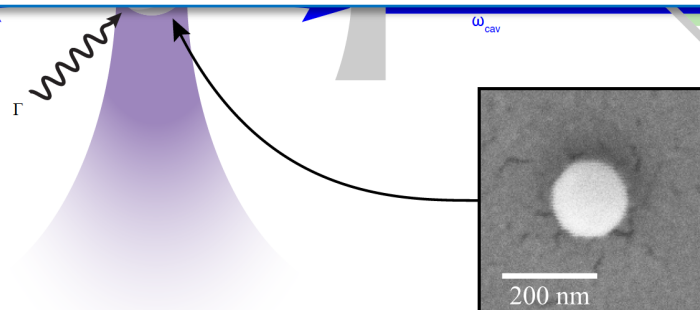
$$\bar{n} = \left(\frac{\kappa}{4\Omega_x}\right)^2 + \frac{\gamma_g \frac{k_B T_{bath}}{\hbar\Omega_x} + \Gamma_{rec}}{\Gamma_- - \Gamma_+} + \frac{n_{phot}}{\kappa} S_\varphi(\Omega_x)$$

0.026

2×10^5 @ 0.06 mbar

$< 10^{-3}$

0.5 @ 1e-6 mbar



$\longleftrightarrow \omega_x \approx 2\pi \times 305$ kHz

$\updownarrow \omega_z \approx 2\pi \times 80$ kHz

$\swarrow \omega_y \approx 2\pi \times 275$ kHz

$\kappa = 2\pi \times 193$ kHz

$p = 1e-6$ mbar, $T = 300$ K

$n_x < 0.5$ (ground state probability $> 2/3$)

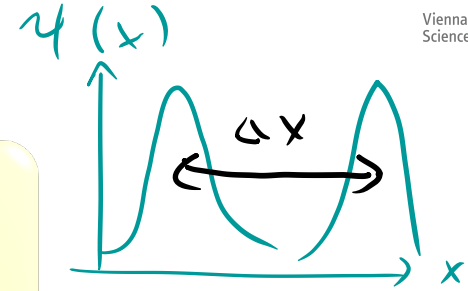
Center-of-mass $T_c = 12$ uK; environment $T_e > 300$ K

$g_x = 2\pi \times 71$ kHz, Cooperativity $C = 5$

$$\dot{\rho} = -\frac{i}{\hbar}[H, \rho] + \mathcal{L}[\rho]$$

Master equation approach

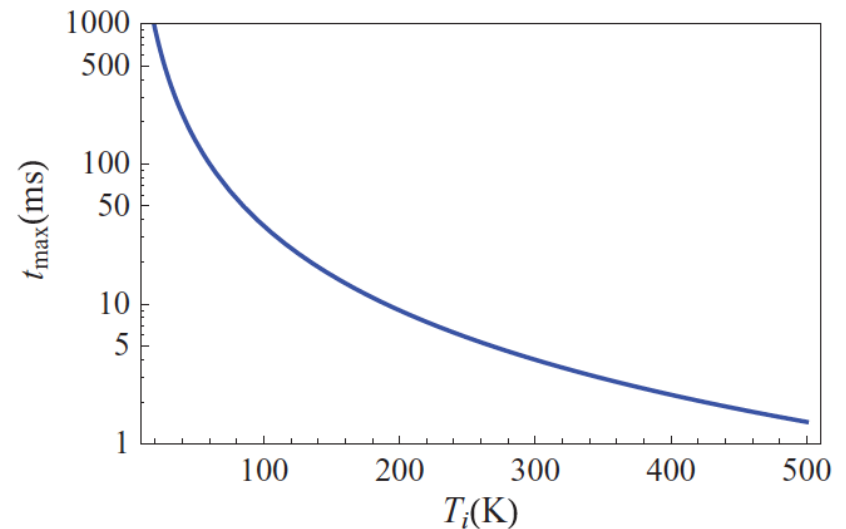
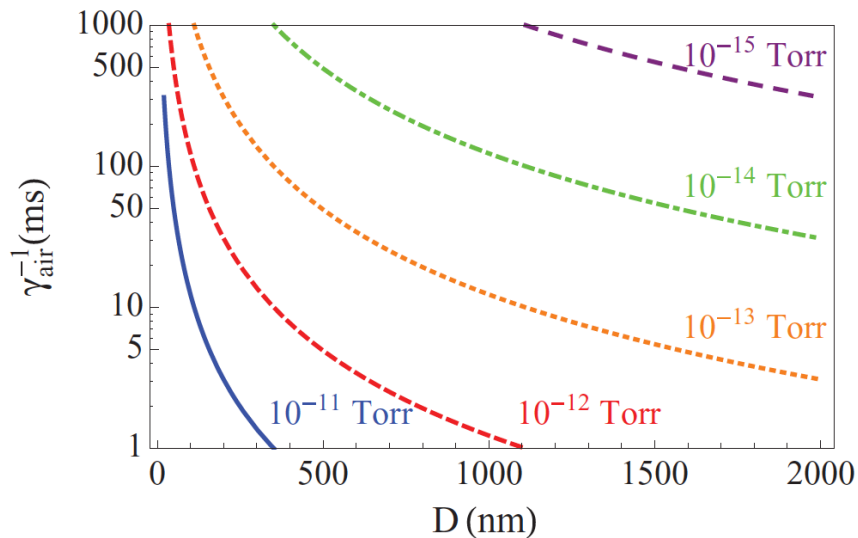
See also
O. Romero-Isart et al.,
PRL 107, 020405 (2011)
O. Romero-Isart, PRA
84, 052121 (2011)



Example: a free nanoparticle

Decoherence due to **gas scattering** on a
glass sphere (Romero-Isart 2011)

Decoherence due to **blackbody absorption**
(50 nm sphere)



In our case ($p=1e-6$ mbar; $T_e > 300K$): $\Gamma_{\text{gas}}/2\pi = 15\text{kHz}$, $\Gamma_{\text{rec}}/2\pi = 6\text{kHz}$

$$\dot{\rho} =$$

Photon Recoil & Gas Scattering limit **in-trap coherence time** to $< 8\mu\text{s}$
(**15 coherent oscillations**)

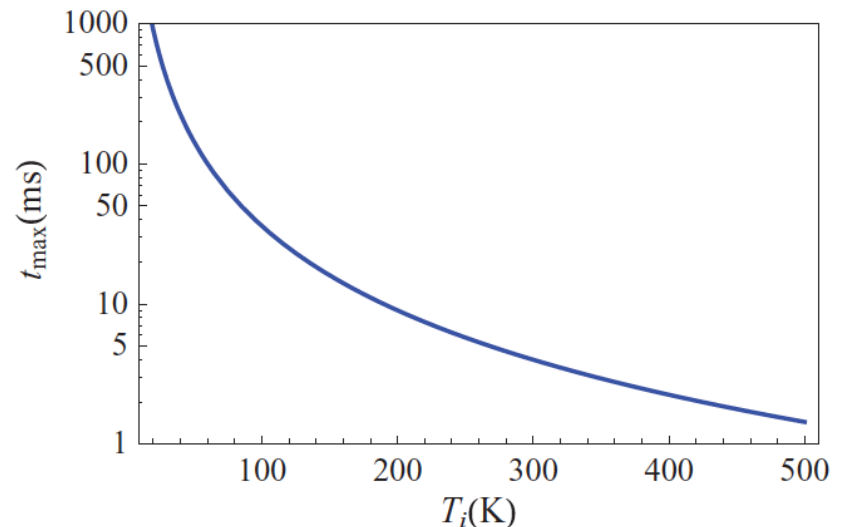
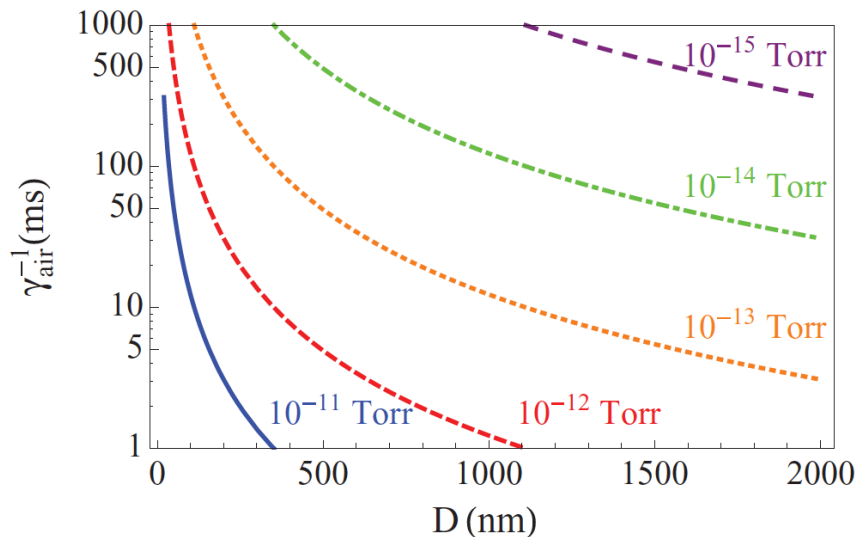
Gas Scattering limits **free-fall coherence time** to $< 2\mu\text{s}$
(**wavepacket expansion by factor of 3**: $3\text{pm} \rightarrow 10\text{pm}$)

Wavepacket size $>$ particle size will require $p < 1e-11\text{mbar}$ and $T_e < 130K$

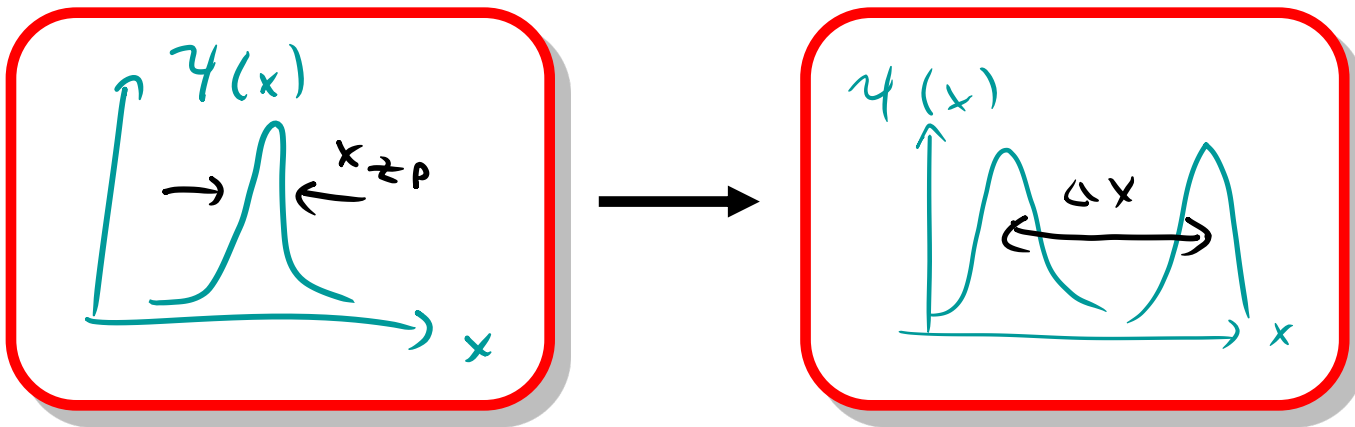
Examp

Decoherence
glass sphere (Romero-Isart 2011)

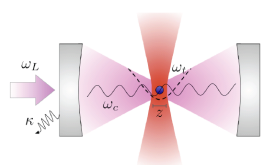
Decoherence
(50 nm sphere)



Towards „large“ quantum superposition states

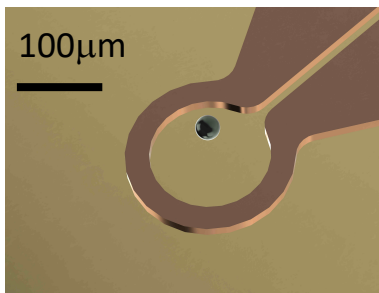
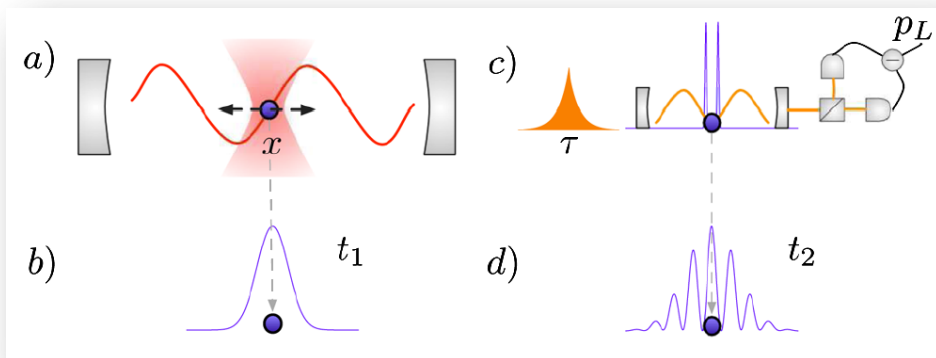


Free-fall + quantum measurement



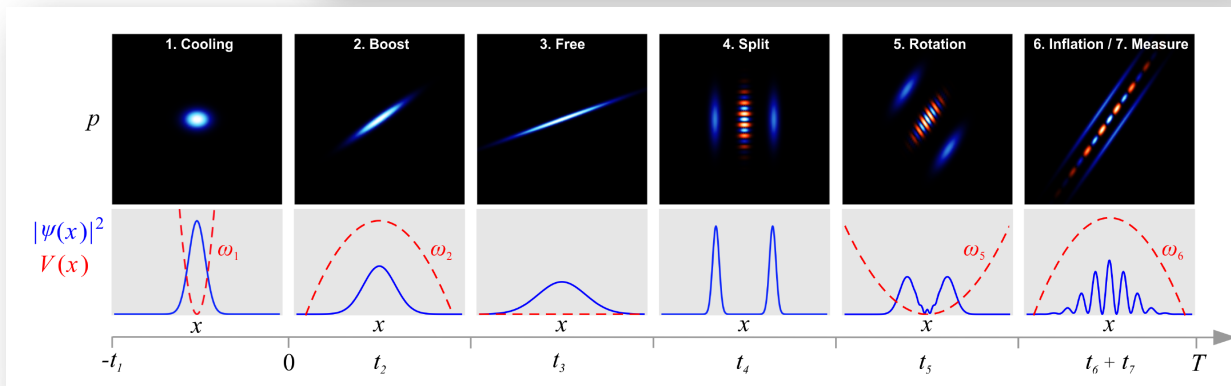
Optical levitation

(Romero-Isart 2011)
PRL107, 020405

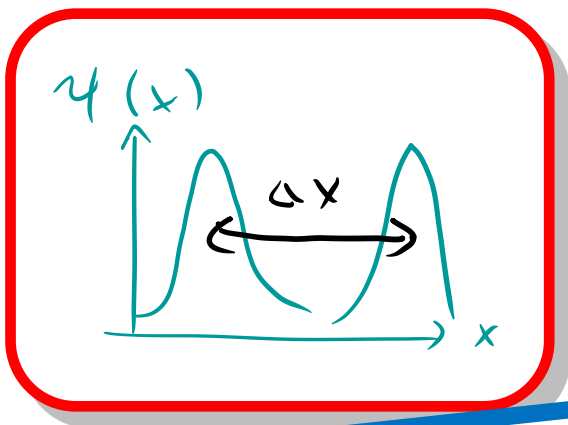
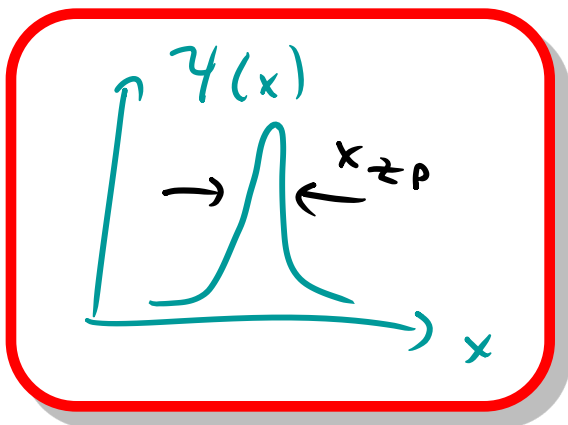


Superconducting levitation

(Pino 2016) arxiv: 1603.01553

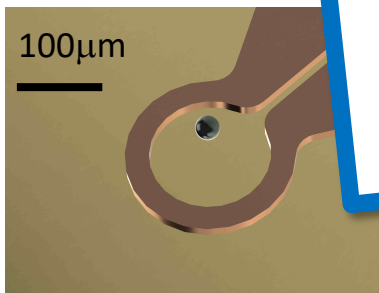
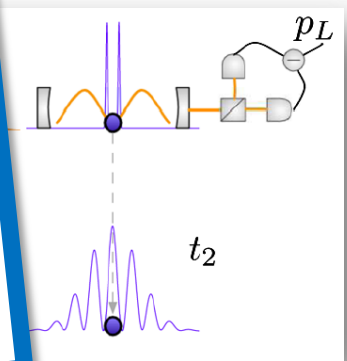
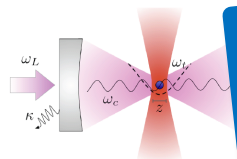


Towards „large“ quantum superposition states



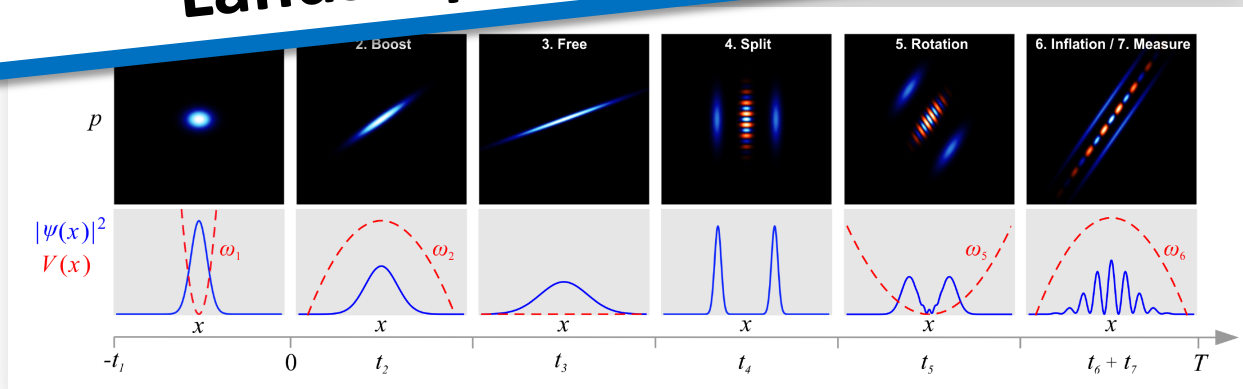
Free-fall + quantum measurement

**Next Step:
Controlling the Potential Landscape**



Superconducting levitation

(Pino 2016) arxiv: 1603.01553




 Mario Ciampini
(Kiesel group)

Shaping the potential landscape of optical tweezers

OPTICAL LEVITATION:

$$\hat{H} \ll \underline{d} \cdot \underline{E} = \chi \cdot E^2$$

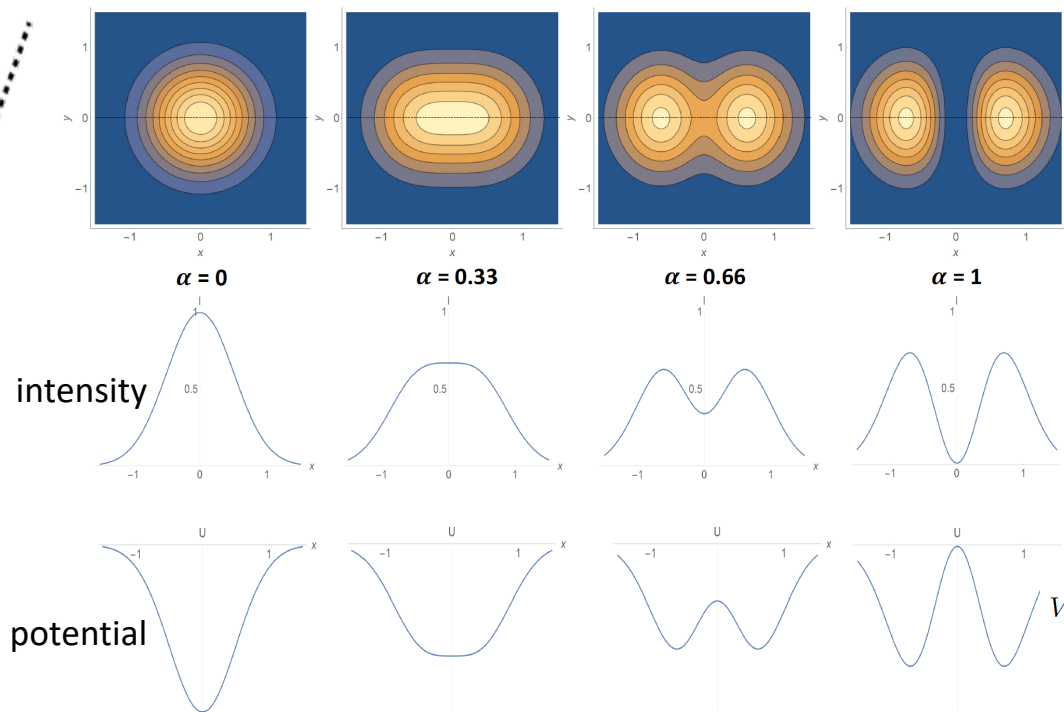
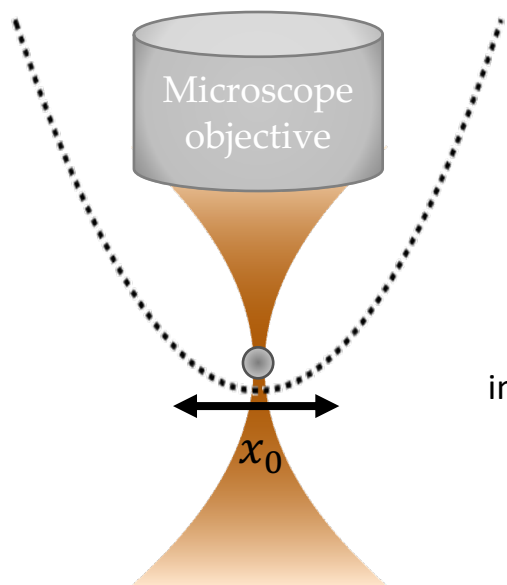
χ : $\text{Re}\{\text{Polarizability}\}$
 E : optical trapping field

↳ beam intensity

$$\rightarrow \text{GRADIENT FORCE } F \propto (\nabla E^2) \cdot \chi$$

Gaussian TEM₀₀ provides 3D harmonic trap (to first order)

Superposition of TEM₀₀ with TEM₀₁ provides control over potential landscape, e.g. from harmonic to repulsive

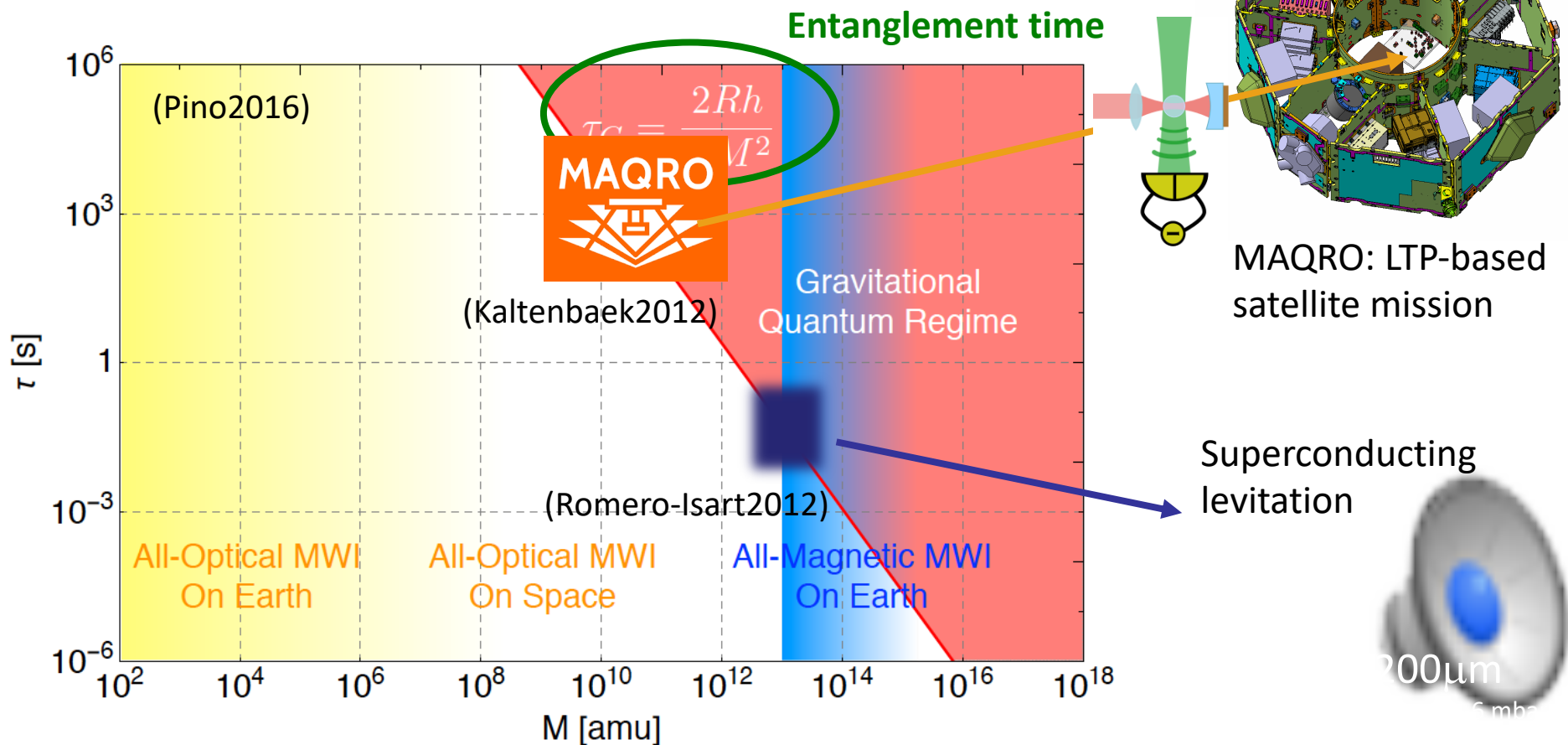
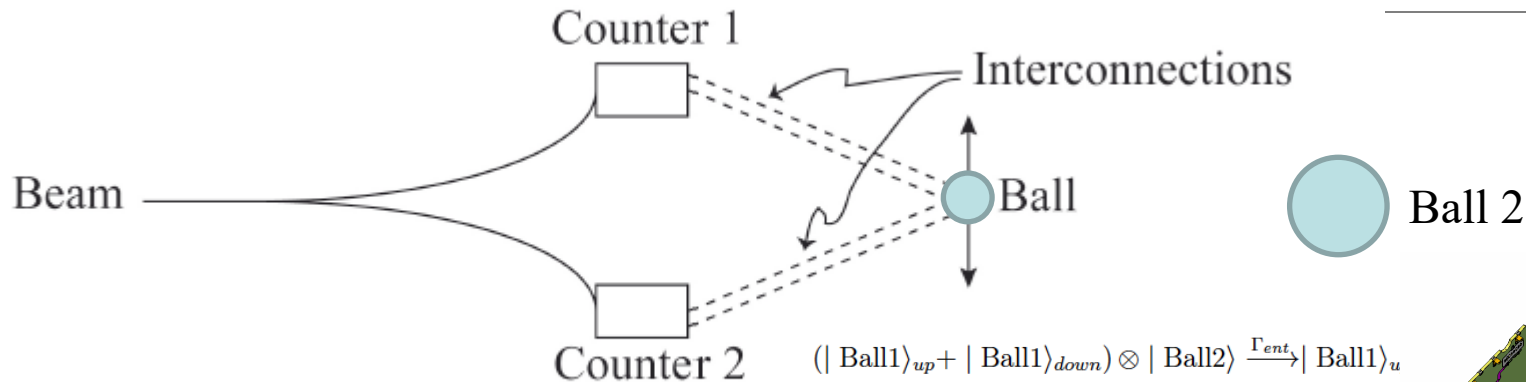


$$\alpha \text{TEM}_{00}(H) + \beta \text{TEM}_{01}(V)$$

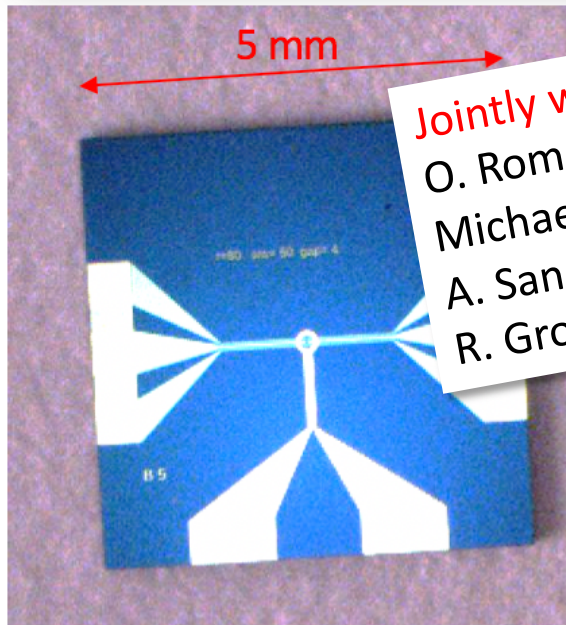
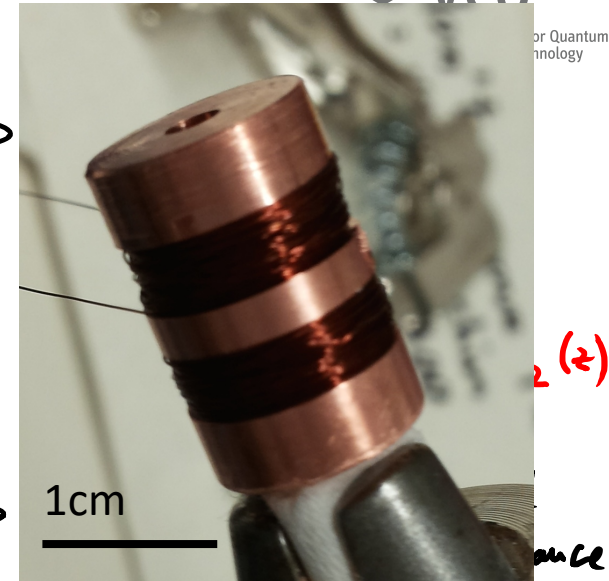
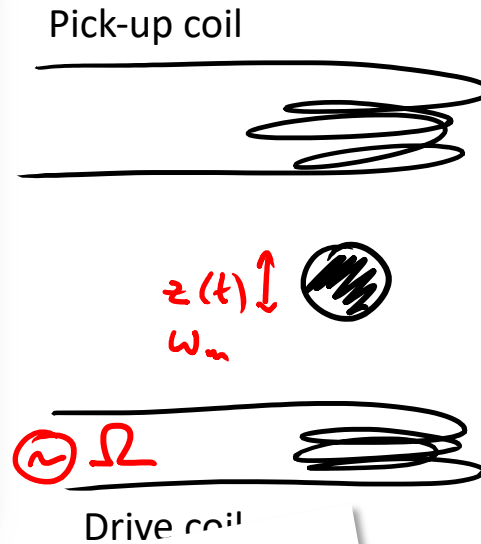
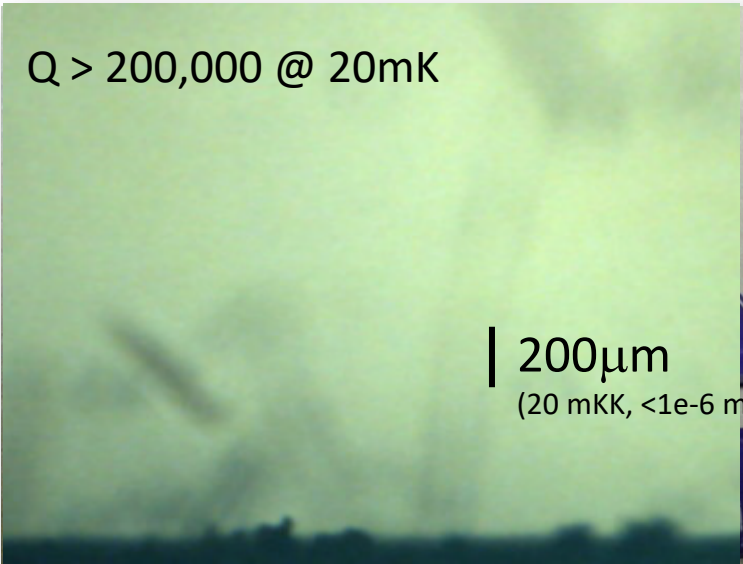


$$V(x, t) = - \left[\alpha + \frac{\beta}{2} x^2 \right] \exp\left(\frac{-x^2}{2\omega_0^2}\right)$$

An ultimate experiment? Entanglement by gravity...

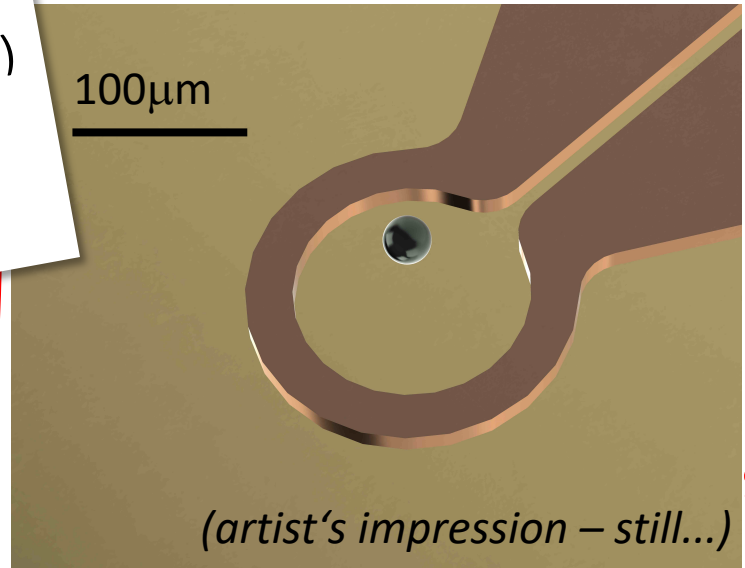


Magnetically trapped superconductors as mechanical resonators



Jointly with:
 O. Romero-Isart, G. Kirchmair (IQOQI)
 Michael Trupke (U Vienna)
 A. Sanchez (UA Barcelona)
 R. Gross, H. Huebel (WMI Munich)

Magnetic levitation in anti-Helmholtz coil configuration
 Trap frequencies ~ 1 kHz
 $T = 20$ mK, $p = 1\text{e-}8$ mbar



J. Hofer, S. Miniberger, M. Trupke

Optical levitation in cavities (with **Kiesel Group**, V. Vuletic)

Uros Delic

David Grass (@Duke)

Constanze Bach

Yuriy Coroli

Jelena Cvijan

Kahan Dare

Lorenzo Magrini

Manuel Reisenbauer

Superconducting levitation (with R. Gross, O. Romero-Isart, M. Trupke)

Josh Slater (@ Delft)

Stefan Minniberger

Milan Gemaljevic

Joachim Hofer

Quantum foundations and the gravity-quantum interface (with C. Brukner, B. Dakic, R. Wald, A. Zeilinger)

Alessio Belenchia (@ Belfast)

Lukas Neumeier

Fatemeh Bibek

Philipp Köhler

Potential landscape shaping & optimal control (with **Kiesel Group**, A. Kugi, M. Ritsch-Marte)

Mario Ciampini

Maxime Debiossac

Stefan Lindner

Tobias Wenzl

Qianze Zhu

Quantum information interfaces (with K. Hammerer, S. Gröblacher, O. Painter, R. Schnabel, J. Eisert)

Sungkun Hong (@ KAIST)

Ralf Riedinger (@ Harvard)

Witlef Wieczorek (@ Chalmers)

Claus Gärtner

Corentin Gut

Klemens Winkler

Precision measurements of gravity

Tobias Westphal

Mathias Dragosits

Hans Hepach

Jeremias Pfaff

Low-noise coatings & microfab

Garrett Cole @ CMS

Quantum Optomechanics at ultra-low temperatures

HBT of single phonons

Hong, Riedinger et al., *Science* **358**, 203 (2017)

Entangled mechanical oscillators

Riedinger, Wallucks, et al. *Nature* **556**, 473 (2018)

Marinković et al., *Phys. Rev. Lett.* **121**, 220404 (2018)

The gravity-quantum interface

Quantum Superposition of Massive Objects and the Quantization of Gravity

Belenchia et al., *Phys. Rev. D* **98**, 126009 (2018)

Levitating dielectrics and superconductors

Near-field coupling to a photonic crystal cavity

Magrini et al., *Optica* **5**, 1597 (2018)

Dispersive cavity cooling in UHV

Delic et al., arXiv:1902.06605 (2019)

Cavity cooling via coherent scattering

Delic et al., *Phys. Rev. Lett.* **122**, 123602 (2019)

Thanks!

