Laplacian Growth

Singular limit of dispersive waves and related problems in quantum and classical physics

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Laplacian growth is an evolution of a plane interface between two immiscible phases, driven by a gradient of a harmonic field



$$\vec{v}_n = \vec{\nabla} p$$

$$\Delta p = 0$$

 $p \to \log |z|, \quad |z| \to \infty$ p = 0 on the boundary

normal velocity = gradient of harmonic field



• Stochastic Growth

Zabrodin

Agam

Bettelheim

Krichever

Mineev-Weinstein

Teodorescu

Hele-Shaw cell



Oil (exterior)-incompressible liquid with high viscosity

Water (interior) - incompressible liquid with high viscosity

Laplacian growth;

diffusion driven patterns.

 $v_n = -
abla_n P$ on the interface

$$\Delta P = 0$$
 in oil,

 $P = \sigma imes$ curvature in water

σ -surface tension

Velocity= gradient of a harmonic function

Fingering Instability





Large flux, small surface tension_ Small Flux, large surface tension_



Fingering instability



Flame with no convection



after H. Swinney et al

Finite time singularities



Hypotrochoid grows until it reaches a critical point.

Reformulation:

Exterior harmonic moments

$$t_k = -\frac{1}{k} \int_{exterior} (x + iy)^{-k} dx dy$$

= Conserved

Find an evolution of a domain domain, which area grows, while the moments are conserved

 t_k - are coordinates of the moduli space of the droplet

- Fingering Instability
- Finite time cusp-like singularities
- Necessity of a regularization of short distances

Interpretation:

pr

essure -
$$p = \log |w(z)|$$
 conformal map

velocity - $v = \nabla p = |w'(z)|$ conformal measure

A probability of a Brownian mover to arrive and join the aggregate



While joining the aggregate a mover becomes a *fermion*.



Evolution of conformal map

• Evolution of Riemann's surfaces:



Degenerate Riemann surface: all cycles are nulled, except one



Diffusion limited aggregation







Fluid dynamics

Turbulent Regime, DLA

Arriving particles have a finite size: \hbar

DLA - Diffusion limited aggregation

Noise - a mechanism of stabilization

- Saffman
- Taylor
- Kruskal
- Richardson
- Kadanoff
- Shraiman
- Hastings
- Levitov

Recent developments:

- Link to Random (non-Hermitian) Matrix Ensembles particles are eigenvalues of a random matrix;
- Link to QHE and Correlation functions in quantum wires: water domain= Quantum droplet on the LLL, or a Fermi surface in the phase space;
- Link to a Boundary Conformal Field Theory;
- Link to Integrable system Toda lattice hierarchy (singular limit of dispersive waves)
- Link to problems of 2D Quantum gravity

Aharonov-Bohm Effect on LLL

Particles on LLL in a uniform magnetic field form a circular droplet



 $\Psi(z_1, ..., z_N) = \prod (z_n - z_m) e^{-\sum_n \frac{|z_n|^2}{2\ell^2}}$ m < n



Three interpretations:

• 2D-Electrons on the first Landau level (electronic droplet);

 $\Psi(z_1, ..., z_N) = \prod_{m < n} (z_n - z_m) e^{-\sum_n \frac{|z_n|^2}{2\ell^2}}$ • Phase space (coherent state reprs) of ID Electrons in parabolic potential;

 $\Psi(z_1,...,z_N)\,$ - Wigner (or coherent state repr.) $p^2+E^2=N\hbar$

•Eigenvalues distribution of general Random Matrix ensemble

$$M = U^{-1} \operatorname{diag}(z_1, \ldots, z_N) U$$

$$e^{-\frac{1}{\hbar}\operatorname{Tr} MM^{\dagger}} dM \sim \prod_{m < n} |z_n - z_m|^2 e^{-\frac{1}{\hbar}\sum_n |z_n|^2}$$

The circular shape is unstable under gradients of magnetic field



Magnetic field = $B = -\partial_z \partial_{\bar{z}} (|z|^2 - V(z) - \overline{V(z)})$

$$\Psi(z_1, \cdots, z_N) = \frac{1}{\sqrt{N!\tau_N}} \Delta(z) e^{-(\sum_n \frac{1}{2\ell^2} |z_n|^2 - V(z_n))}.$$

The saddle point equation (4) is transformed accordingly

$$\sum_{m \neq n}^{N} \frac{2\ell^2}{z_n - z_m} = \bar{z}_n - 2\ell^2 \frac{\partial}{\partial z} V(z).$$

Semiclassical approximation: $\hbar \to 0$, $N \to \infty$, $N\hbar =$ fixed

The circular shape is unstable under gradients of magnetic field





Semiclassical limit of the Growth: N→ N+1 provides the same hydrodynamics as Laplacian growth.

Quantum limit provides stochastic regularization

- Quantization a way of regularization of classical systems at a singular regime:
- Quantization introduces a new scale

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Cusp-like singularities



- Hydrodynamics of Dispersive waves;
- 2D-gravity
- Random matrices (edge of the spectrum)

Hermitian matrix ensemble:

eigenvalues are real, distributed in an interval of the real axis



Figure 2.2: Typical spectrum of a matrix. In the large N limit, the spectrum is well approximated by a continuous density of levels $\rho(\lambda)$.

Bulk properties are universal (do not depend on potential V(M))

Edge properties are also universal but do depend on potential



e 2.6: The density of eigenvalues has square-root singularities at the edges. An ement of the edge's vicinity shows a universal tail related to the Airy function.



Non-At Gaussian potential power law singularities at the edges. An enlargement of the edge's vicinity shows a universal tail related to some integrable hierarchy of differential equations.

Edge singularities are characterized by p, q

Quantization can be seen as a noise

Noise is negligible in the semiclassical regime (a smooth droplet)

Important in the turbulent (singular)
 regime, where semiclassics fails.

Singular semiclassical expansion at the cusps

- Laplacian growth=Witham hierarchy of soliton equations
- Dispersive regularization of nonlinear waves

Connection to integrable systems

How does the shape of the domain change with respect to a change of harmonic moments t_k ?

d (dispersionless) Toda hierarchy

How does the wave function change with respect to a change of harmonic moments t_k and \hbar ?

Toda hierarchy

- Witham hierarchy = a semiclassical limit of soliton equations
- Example KdV u(x,t):

 $|\mathbf{u}_t - \boldsymbol{u} \cdot \boldsymbol{u}_x + \boldsymbol{h} \cdot \boldsymbol{u}_{xxx} = 0$

• Periodic (soliton-like) solution: waves in shallow water;

• Non-periodic solutions;







Overhang - a result of ill- approximation Oscillations is a dispersive regularization

$$u_t - u \cdot u_x + \hbar \cdot u_{xxx} = 0 \quad \Rightarrow \quad u_t - u \cdot u_x = 0$$

Laplacian Growth is = dispersionless Toda hierarchy

Quantized Laplacian growth (QHE, or Random Matrix Models, etc.) = Toda Integrable Hierarchy



Fine structure of a cusp $P(z) = |\Psi(z)|^2$ Painleve I equation $\left(-\frac{d^2}{dt^2} + u(t)\right)\Psi(z) = z\Psi(z)$ u(t) - a coordinate of a tip

Cusp (2,3):

$$\hbar \cdot u_{tt} - u^2 = t - t_c$$

Universal character of the fine structure of singularities

u(x,t)

• Singularities are classified by two integers (p,q) according to Kac table of CFT

•The simplest cusp (2,3) - Painleve I equation

•Higher order singularities - higher order Painleve like equations

Asymmetries and tip splitting



Set of operators perturbing singularities Set of exponents - gravitational descendants

$$\Delta_{p,q}^{m} = \frac{q - p + m}{p + q - 1}, \quad m = 1, \dots, q < p$$

Are fractal exponents of DLA among this set?