

*Danmarks Grundforskningsfond - Quantum Optics Center*



**Entanglement of simple  
macroscopic objects**

**Niels Bohr Institute - Copenhagen**

*Quantum mechanical wonders  
(second wave)*



**Quantum objects**

**cannot be measured**

**cannot be copied**

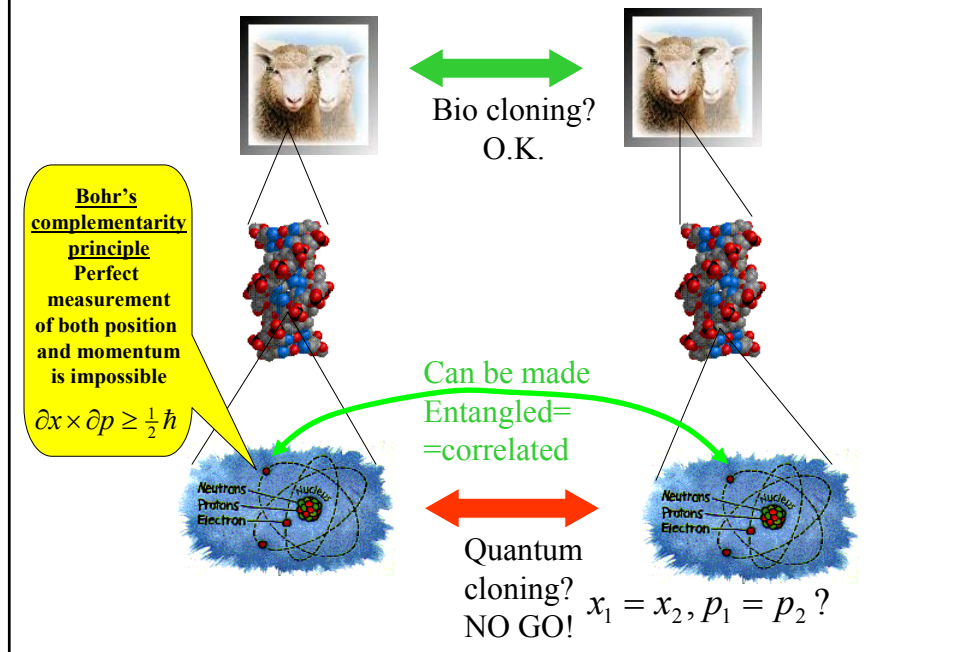
**exist in superposition**

**and entangled states**

**Quantum Information Science**

- Communications with  
absolute security
- Computing with unprecedented speed
- Teleportation of objects (or at least of  
their quantum states)
- Quantum memory

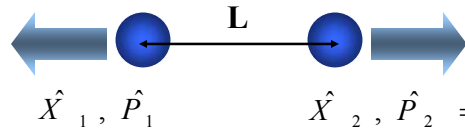
Nothing in this world can be perfectly measured, or perfectly cloned...



*"Entanglement is ... perhaps the most fundamental issue in quantum mechanics" – Erwin Schrödinger*

• *Einstein-Podolsky-Rosen example 1935*

2 particles entangled in position/momentum



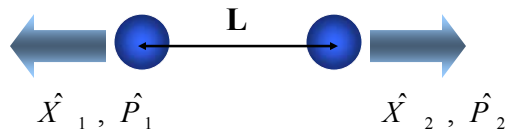
$$\hat{X}_1 - \hat{X}_2 = L \quad \hat{P}_1 + \hat{P}_2 = 0$$

*Simon (2000); Duan, Giedke, Cirac, Zoller (2000)*

Necessary and sufficient condition for entanglement

$$\delta ( X_1 - X_2 )^2 + \delta ( P_1 + P_2 )^2 < 2 \hbar$$

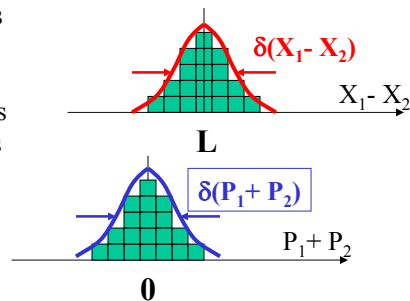
2 particles entangled in position/momentum



**What does it mean in practice?**

$$\delta ( X_1 - X_2 )^2 + \delta ( P_1 + P_2 )^2 < 2 \hbar$$

1. Prepare many identical pairs of particles
2. Measure  $X_1 - X_2$  on some of those pairs
3. Measure  $P_1 + P_2$  on others
4. Plot statistical distributions of the results
5. Measure the width of these distributions



## Why does it make sense?

### Two independent particles

$$\hat{X}_1, \hat{P}_1$$

$$\hat{X}_2, \hat{P}_2$$

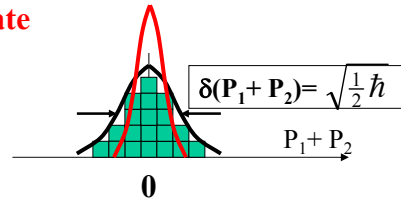
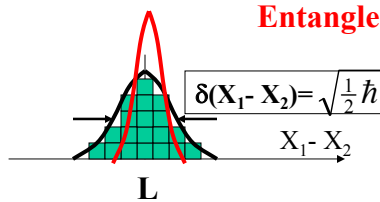
$$\delta x \times \delta p \geq \frac{1}{2} \hbar$$

Minimal symmetric  
uncertainties

$$\delta x_{1,2}^2 = \delta p_{1,2}^2 = \frac{1}{2} \hbar$$

$$\delta (X_1 - X_2)^2 + \delta (P_1 + P_2)^2 \leq 2 \hbar$$

### Entangled state



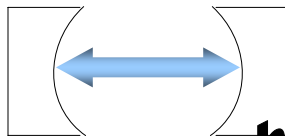
## Position – momentum uncertainty

$$\delta x \times \delta p \geq \frac{1}{2} \hbar \approx 10^{-34} \text{ J sec}$$

### Macroscopic object – a mirror

The best optical interferometry  
measurement:

$$\delta x \approx 10^{-12} \text{ m}$$



Assume  
 $M = 1 \text{ mg}$

**tough!**  $\delta p \approx 10^{-19} \text{ m / sec}$

## Entanglement - spins

2 quantum coins



2 spins

( spin “up” or spin “down”)



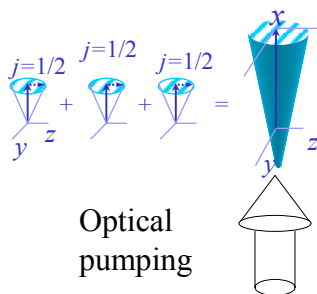
Entangled state:

$$a |00\rangle + b |11\rangle$$

Entangled state, many qubits:



Ensemble of  $N$  atomic spins  
and more about  $\sqrt{N}$



Optical  
pumping

$$J_x = \frac{1}{2}N$$

$$-\frac{1}{2} \quad \frac{1}{2}$$

Along x: all tails



What is the spin projection along Y or Z?

$$-\frac{1}{2} \quad \frac{1}{2}$$

Zero?

Not really!

$$\frac{1}{2}\sqrt{N}$$

Ensemble polarized along X, measured along Y or Z



Binomial => Poissonian  
 $N \gg 1$

Zero mean value with random misbalance  
 between heads and tails  $\frac{1}{2}\sqrt{N}$

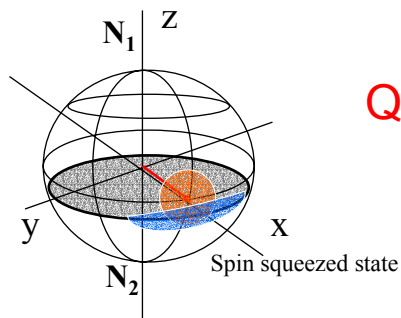
**QM:**  $[\hat{J}_y, \hat{J}_z] = i\hat{J}_x \Rightarrow \delta J_y \delta J_z \geq \frac{1}{2} J_x = \frac{1}{4} N$

**QM: uncorrelated spins =>**  $\delta J_y = \delta J_z = \frac{1}{2}\sqrt{N}$

Alexandrov, Zapassky 1986:  
 "Faraday rotation noise"

Wineland et al 1992: "projection noise"


One atomic ensemble



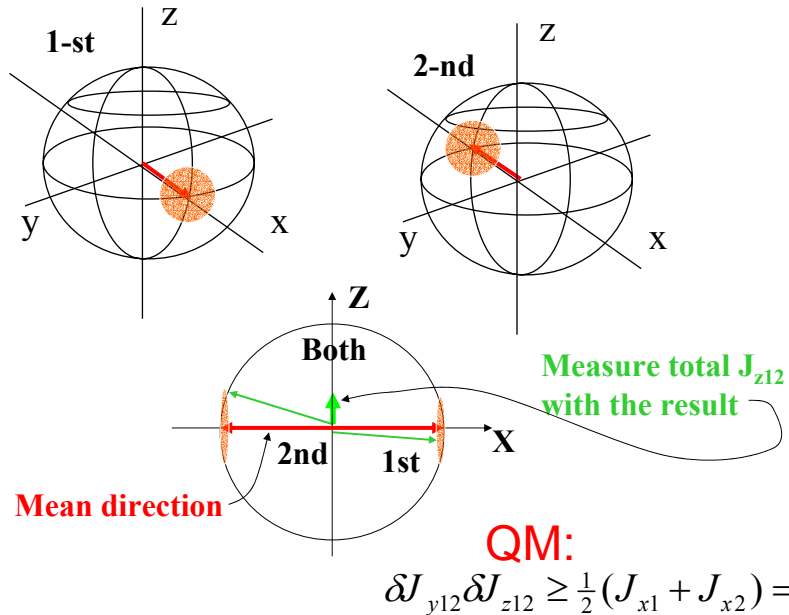
Like P of a particle

**QM:**  $\delta J_y \delta J_z \geq \frac{1}{2} J_x = \frac{1}{4} N$

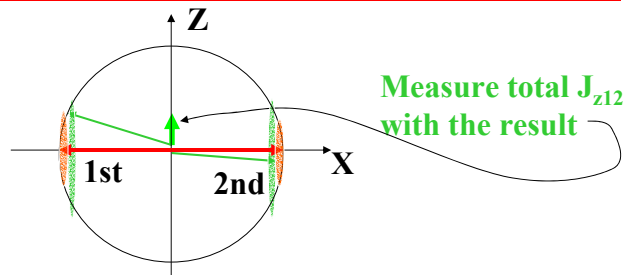
Like X of a particle

$N_1$        $N_2$   
  
 $Z_{\uparrow}$        $Z_{\downarrow}$   
 $Var(N_1 - N_2) < N$

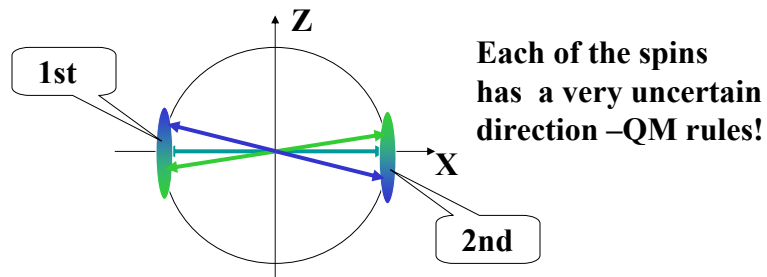
## Two oppositely magnetized atomic ensembles




## Entangled state of two ensembles produced by measurement



If  $J_{z12}$  is known, it can be made zero by spin rotation:  
Now the two spins are collinear !



$\hat{X}_1, \hat{P}_1$  

**Compare**

$$\delta J_y \delta J_z \geq \frac{1}{2} J_x$$

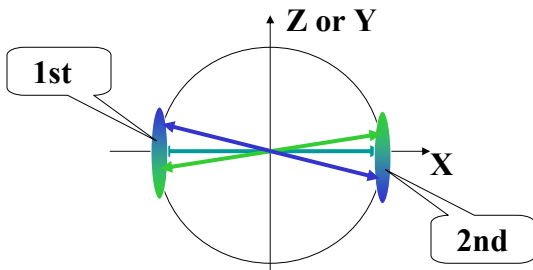
$$\delta x \times \delta p \geq \frac{1}{2} \hbar$$

**If the two macroscopic spins are collinear they must be entangled:**

$$\delta (X_1 - X_2)^2 + \delta (P_1 + P_2)^2 < 2 \hbar$$

**Compare**

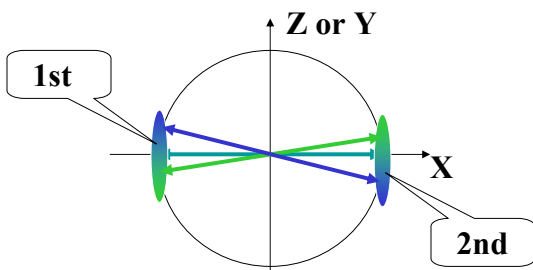
$$\delta (J_{z1} + J_{z2})^2 + \delta (J_{y1} + J_{y2})^2 < 2 J_x$$



**QM:**

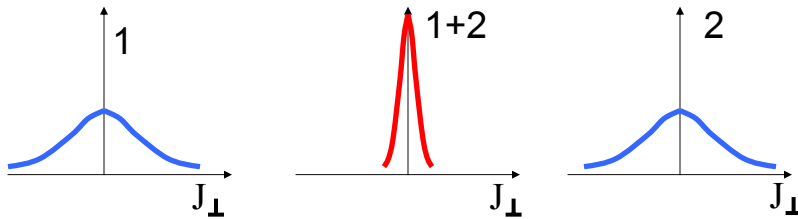
$$\delta J_{y12} \delta J_{z12} \geq \frac{1}{2} (J_{x1} + J_{x2}) = 0$$

**QM allows such a measurement in principle, now all we need is to find the way to perform it...**





Stern-Gerlach projection  
on any axis  $\perp$  to x:



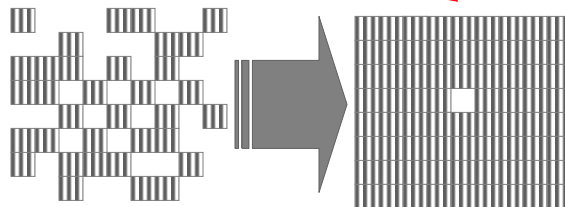
Along y,z: ideally no misbalance between heads and tails of the two ensembles, or, at least, less than random misbalance  $\sqrt{N}$



A secret message  
is encoded in  
one of the grids

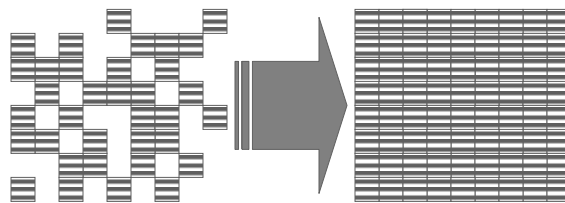
Two Entangled Patterns

Projected  
on vertical  
grid (Y)

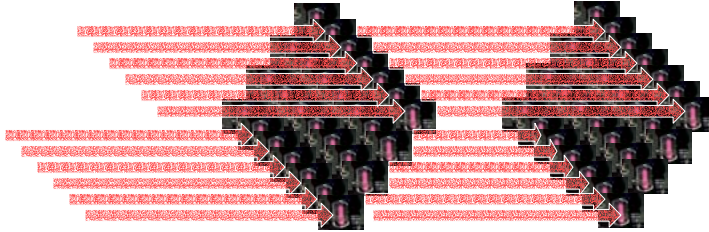


Projection on one grid erases the other grid!

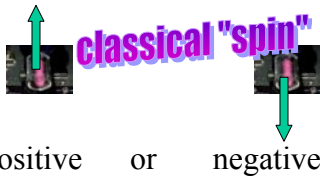
Projected  
on horizontal  
grid (Z)



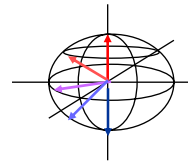
## Communication network: arrays of memory pixels connected with fiber links



Information encoded in magnetization



Quantum spin



## Spin memory with Coherent Spin States

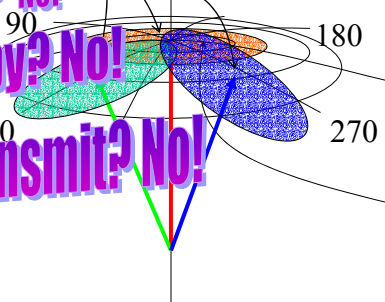
*Quasi-continuous encoding*

*Indistinguishable coherent states*

**Want to read out? No!**

**Want to copy? No!**

**Want to transmit? No!**



$$\delta J_z \delta J_y = 2N$$

- Densely coded states are impossible to read  
but possible to transfer via teleportation-like protocols  
*Teleportation proposal* – Bennett et al '93  
*experiments with light*– Innsbruck'97, Rome'97, Caltech'98

Distributed quantum networking and quantum memory for light can be achieved via atomic teleportation and light-to-atoms state teleportation

**Needed:** distant long-lived  
atomic entangled objects

State-of-the-art experiments:

4 entangled ions in a trap (NIST)

2-3 entangled atoms (ENS, MPQ)

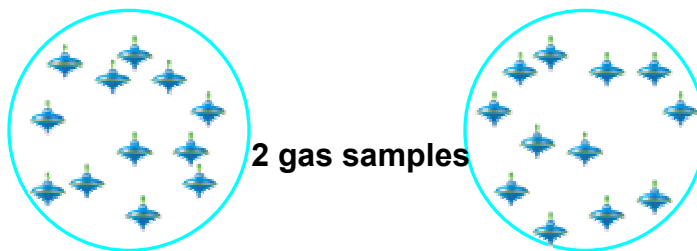
Several atoms in a molecule (NMR)

SQUIDS (?)

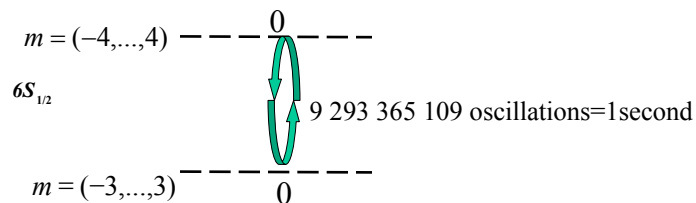


Closely positioned  
objects –  
entanglement via  
localized interactions

$10^{12}$  atoms in each ensemble

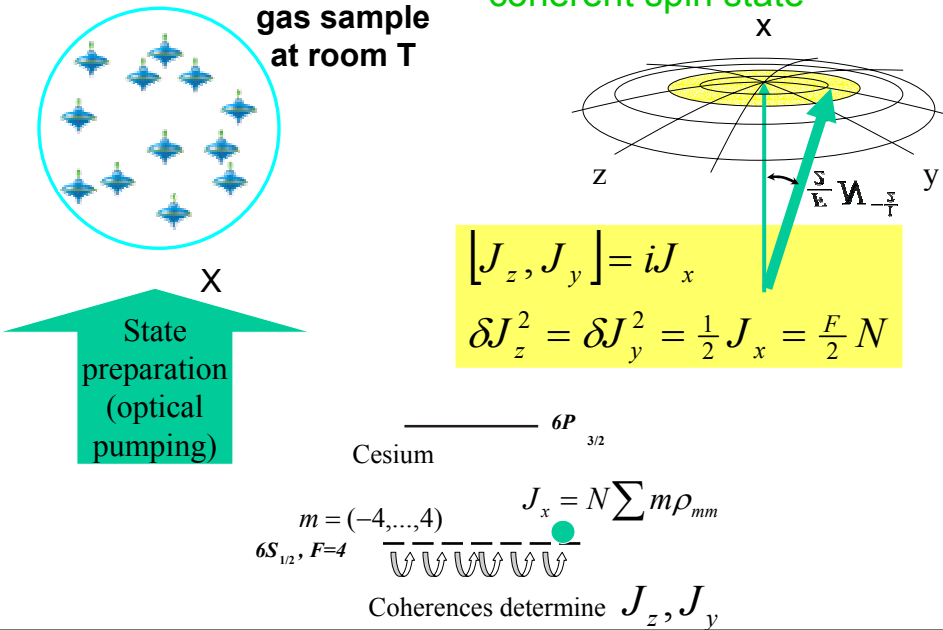


Cesium

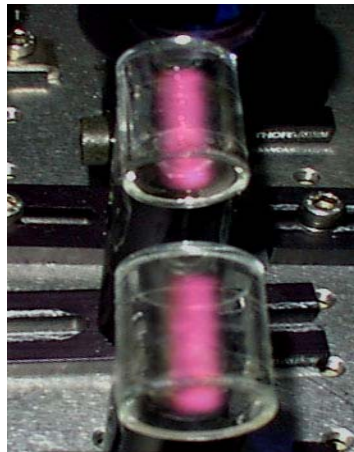


## Macroscopic spin ensemble –

## coherent spin state



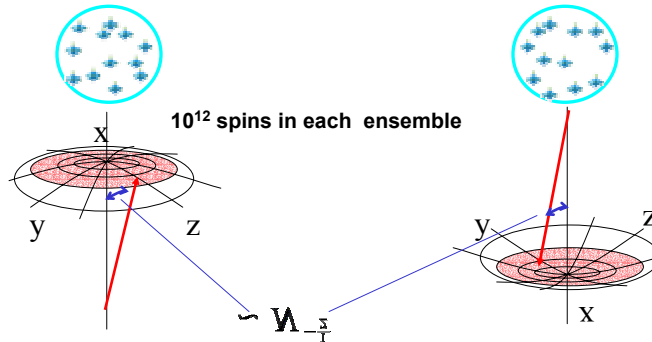
## Entangled state of 2 macroscopic objects



$$\langle (J_{z1} + J_{z2})^2 \rangle + \langle (J_{y1} + J_{y2})^2 \rangle < 2 J_x$$

**Experimental long-lived  
entanglement  
of two macroscopic objects.**

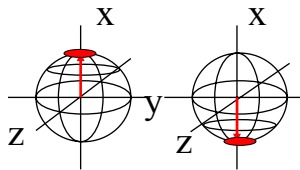
B. Julsgaard, A. Kozhekin,  
and E. S. Polzik,  
**Nature**, 413, 400 (2001).



Spins which are “more parallel” than that  
are entangled

Total Z and Y components of two  
ensembles with **equal and opposite  
macroscopic spins** can be determined  
simultaneously with **arbitrary accuracy**

$$[\hat{J}_{z1} + \hat{J}_{z2}, \hat{J}_{y1} + \hat{J}_{y2}] = i(J_{x1} + J_{x2}) = i(J_x - J_x) = 0$$



Therefore entangled state with

$$\delta(\hat{J}_{z1} + \hat{J}_{z2})^2 + \delta(\hat{J}_{y1} + \hat{J}_{y2})^2 \Rightarrow 0$$

Can be created **by a measurement**

Top view:



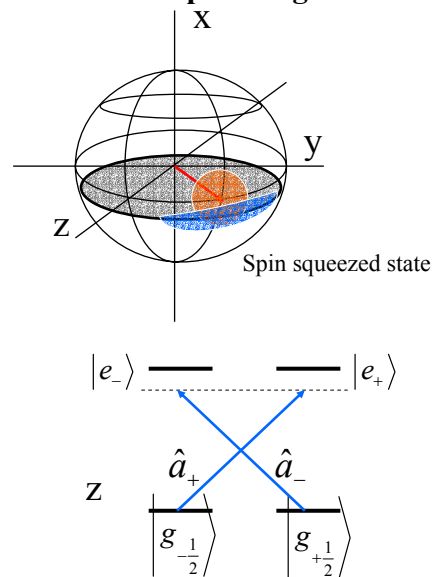
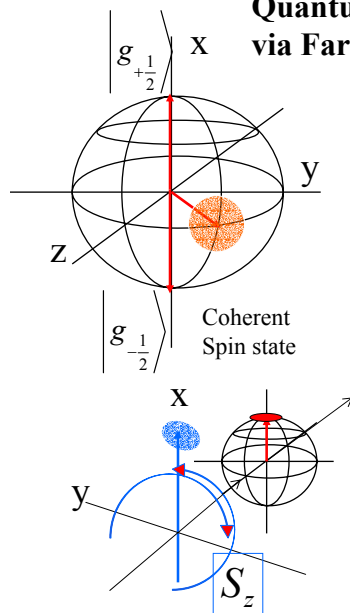
Parallel  
spins must be  
entangled

## How to measure the total spin projections?

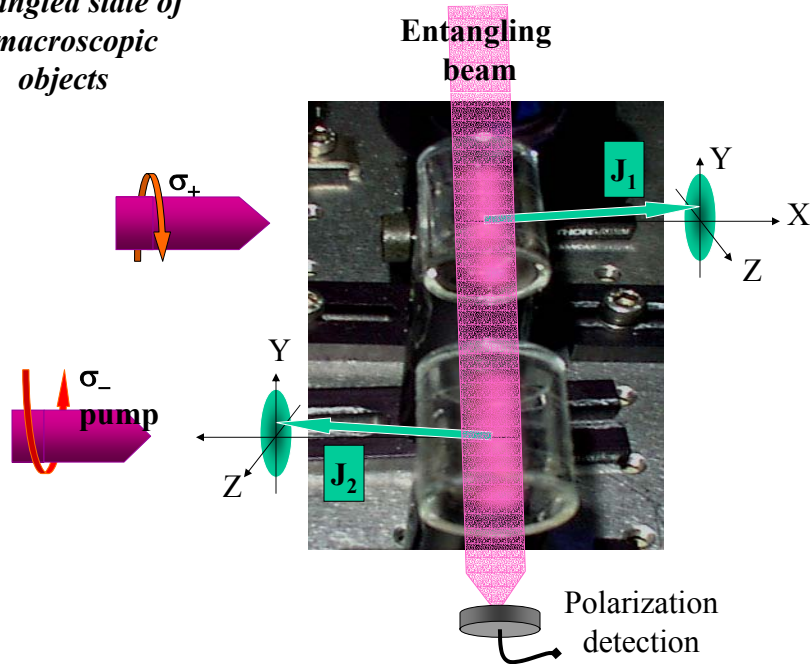
- Send off-resonant light through two atomic samples
- Measure polarization state of light

Duan, Cirac, Zoller, EP 2000

### Quantum non-demolition measurement of spin via Faraday rotation of the probe light



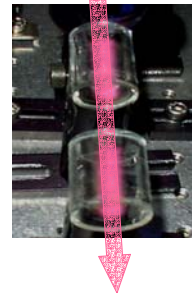
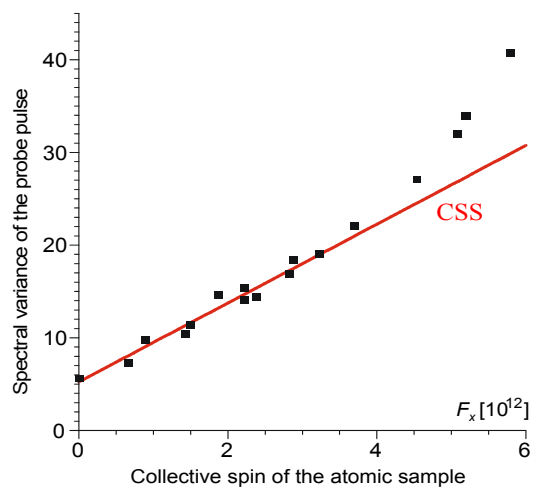
# Entangled state of 2 macroscopic objects



Entanglement criterion:

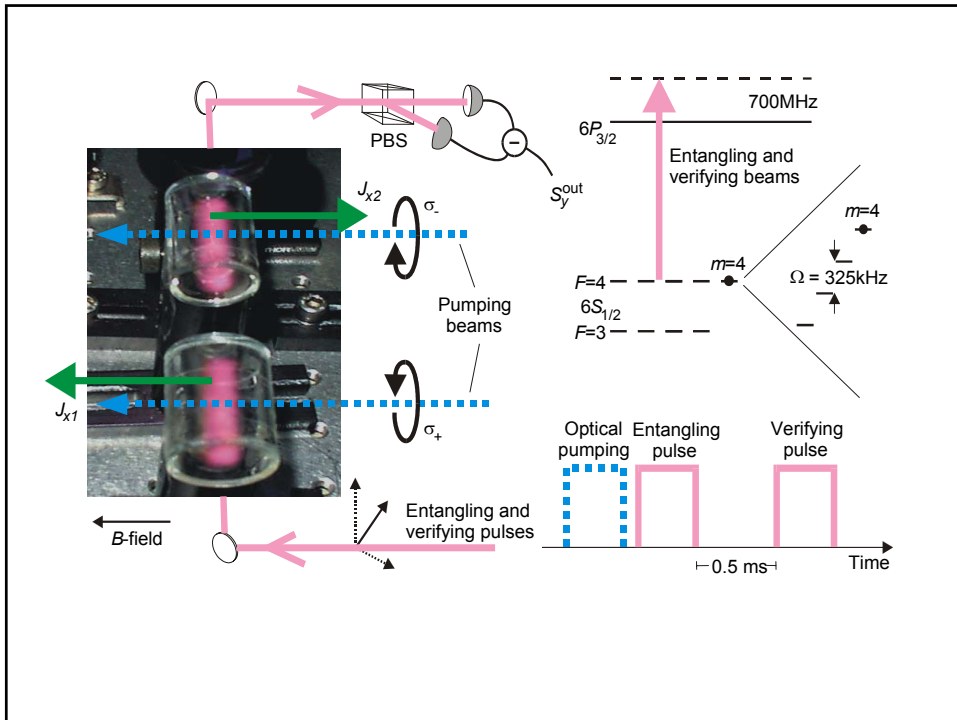
$$\text{Var}(F_{y,1}+F_{y,2}) + \text{Var}(F_{z,1}+F_{z,2}) < 2F_x$$

1) Measure the coherent spin state limit  $2F_x$

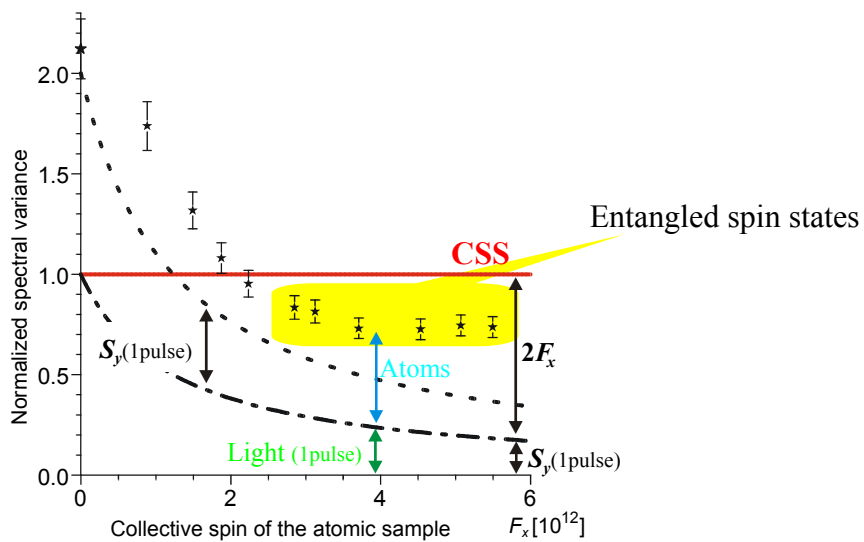


For the coherent spin state:

$$\text{Var}(F_{y,1}+F_{y,2}) + \text{Var}(F_{z,1}+F_{z,2}) = 2F_x$$



## 2) Create entangled state and measure the state variance



*Julsgaard, Kozhekin, EP*

*Nature* **413**, 400 (2001).



## Decoherence issues

- only collective spin states are entangled
- particles are indistinguishable - high symmetry of the system –
  - robustness against losses. This is not a Schrodinger's cat made of  $10^{12}$  atoms!
- no free lunch:
  - limited capabilities compared to ideal maximal entanglement

### Sources of decoherence:

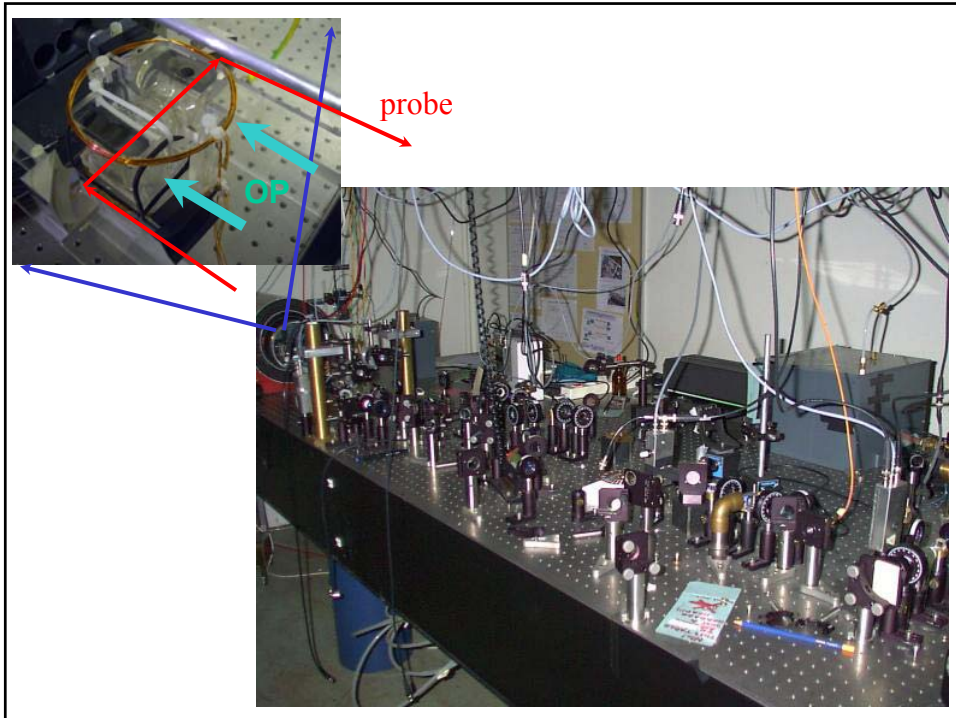
stray magnetic fields	decoherence time 3 milliseconds
collisions	decoherence time 1-2 milliseconds

## Phylosophical issues...

## Phylosophical issues...

**Realism** – two noncommuting spin components  
cannot be measured – therefore do not exist? But can be entangled

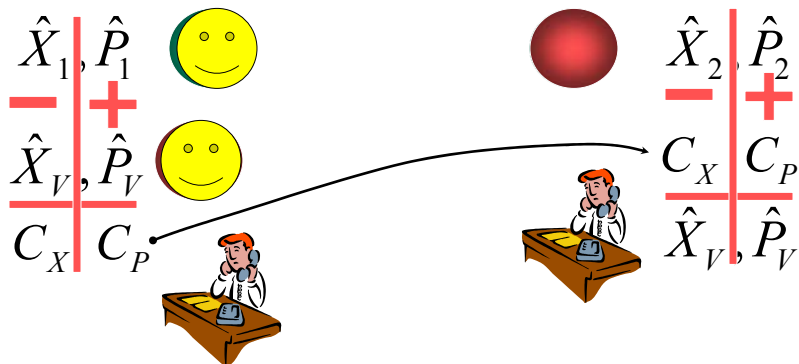
**Non-locality** – two entangled macroscopic objects can be used to  
violate Bell-type inequalities via distillation  
All entangled Gaussian two-mode states are distillable (Giedke et al)



### Teleportation principle (continuous variables)

$$[X, P] = i, [X_1 - X_2, P_1 + P_2] = 0$$

L.Vaidman



Demonstrated experimentally for light variables by Furusawa, Sørensen, Fuchs, Braunstein  
Kimble, Polzik. *Science* 1998

Einstein-Podolsky-Rosen entangled state

$$X_1 - X_2 = 0, P_1 + P_2 = 0$$

Light-to-Atoms Teleportation

$J_{Z1}, J_{Y1}$   
 $- \quad +$   
 $S_{ZV}, S_{YV}$

$J_{Z2}, J_{Y2}$

$$J_y^{out} = J_y^{in} + k S_z$$

$$S_y^{out} = S_y^{in} + k J_z$$

k=1 ↑↓

$\hat{X}_{light}^{out} = \hat{X}_{light}^{in} + \hat{P}_{atoms}^{in}$

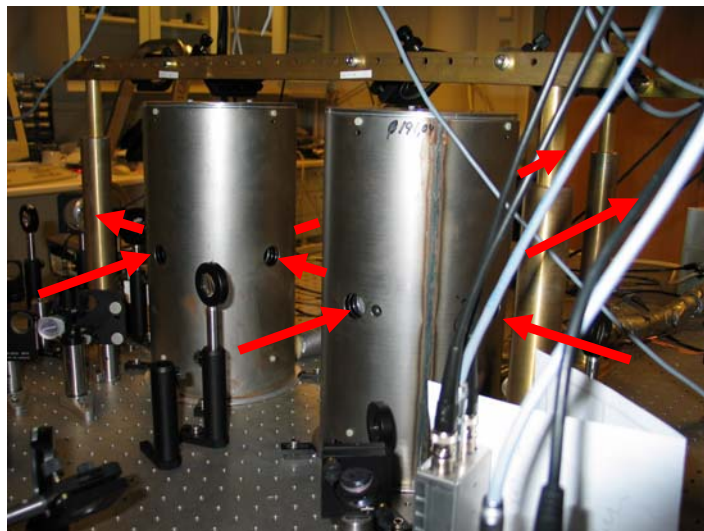
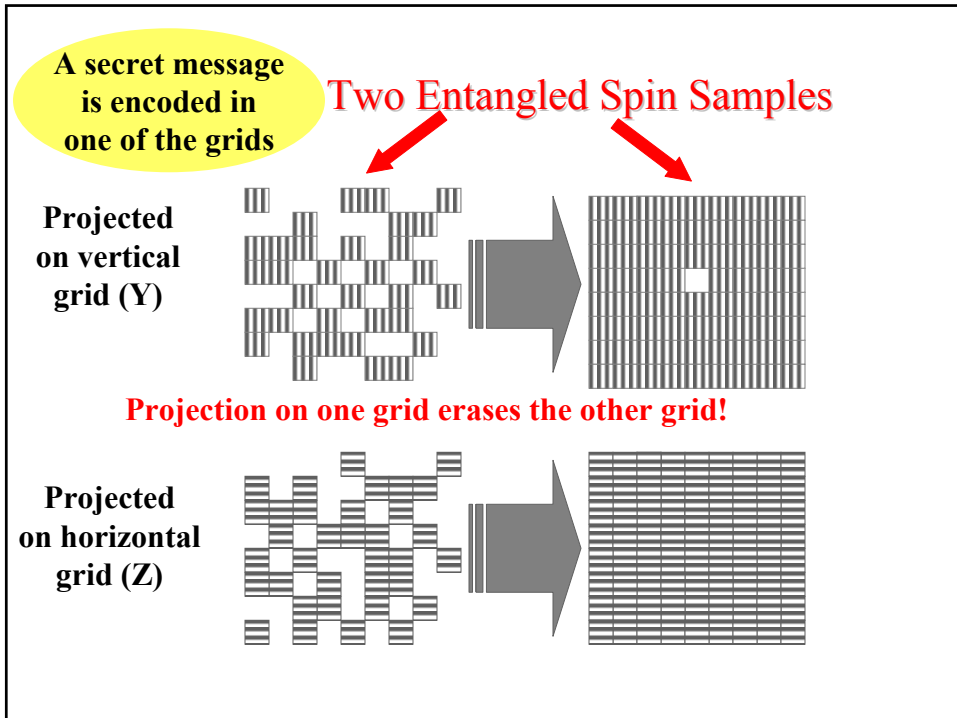
Kuzmich, EP 2000

### Teleportation of an entangled atomic state

- Every measurement changes the single cell spin, BUT does not change the measured sum
- Every pulse measures both y and z components of the sum – entanglement is created

To complete teleportation of entanglement onto cell 1 and cell 4:  
 rotate spin 4 by A+B+C:

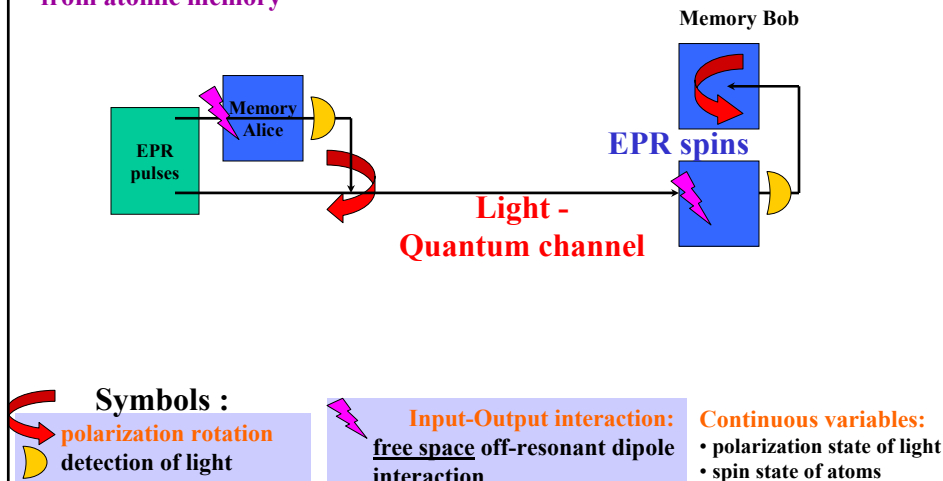
$$\hat{J}_4^{Tel} = \hat{J}_4 - A - B + C = \hat{J}_4 - \hat{J}_1 - \hat{J}_2 - \hat{J}_3 - \hat{J}_4 + \hat{J}_2 + \hat{J}_3 = -\hat{J}_1^{Tel}$$



## Communication networks based on continuous spin variables

**Operation:**  
Storage of light and read-out  
from atomic memory

**Resources:** local entanglement



Scalability – an array of dipole traps or  
solid state implementation – quantum holograms



**Requirements:**

- Long-lived 2 or more level state of a single particle
- The ability to optically pump the system into one of these states
- Dipole coupling to light
- Optically dense medium

**Example:** pencil-shaped Rubidium BEC  
4 microns by 100 microns  
 $10^5$  atoms,  $10^7$  photons per pulse

## Quantum state engineering in Atoms clocks



Cesium atoms

